

# Introduction to S-matrix Bootstrap

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# Outline

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1. Introduction: main idea and axioms
2. The methods: primal and dual approaches
3. Applications: 2D maps, non-invertible symmetries, pions
4. Final Remarks

# Introduction



# The S-matrix Program

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- Motivation: describe strong interactions

- Idea of 1960's program: [e.g. Eden et al 1966, Chew 1966]

Find a unique **non-perturbative S-matrix** by imposing **consistency conditions** coming from general principles ← Bootstrap philosophy

- Modern S-matrix Bootstrap: constrain space of consistent QFTs (*weak and strongly coupled*)

[Paulos, Penedones, Toledo, van Rees, Vieira '16]

→ bound physical parameters

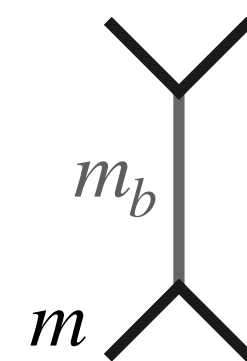
→ chart the landscape of QFTs

# The Axioms (ACU)

## Analyticity (causality)

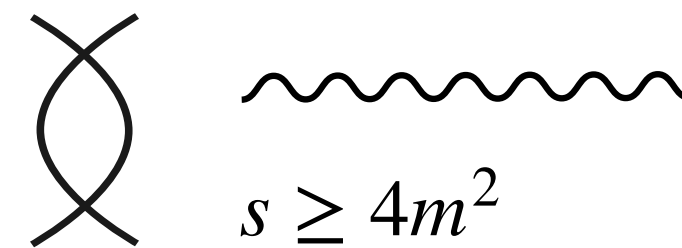
Amplitudes are boundary values of analytic functions of complexified kinematic invariants

**Landau analyticity:** only singularities from on-shell processes



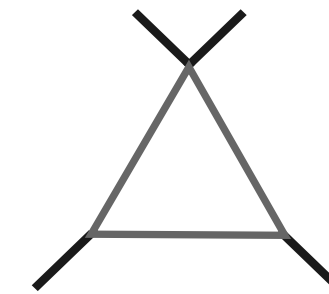
$$\sim \frac{g^2}{s - m_b^2}$$

Bound states



$$s \geq 4m^2$$

Normal thresholds



extra singularities @  $s_*$

*Anomalous* thresholds

\* Only partial proven results for theories with mass gap

## Crossing

Amplitudes in different kinematic channels are boundary values of the same analytic function

Relates process in which we exchange in  $\leftrightarrow$  out particles

\* Beyond 4 particles  $\rightarrow$  inclusive observables [Caron-Huot, Giroux, Hannesdottir, Mizera '23]

## Unitarity

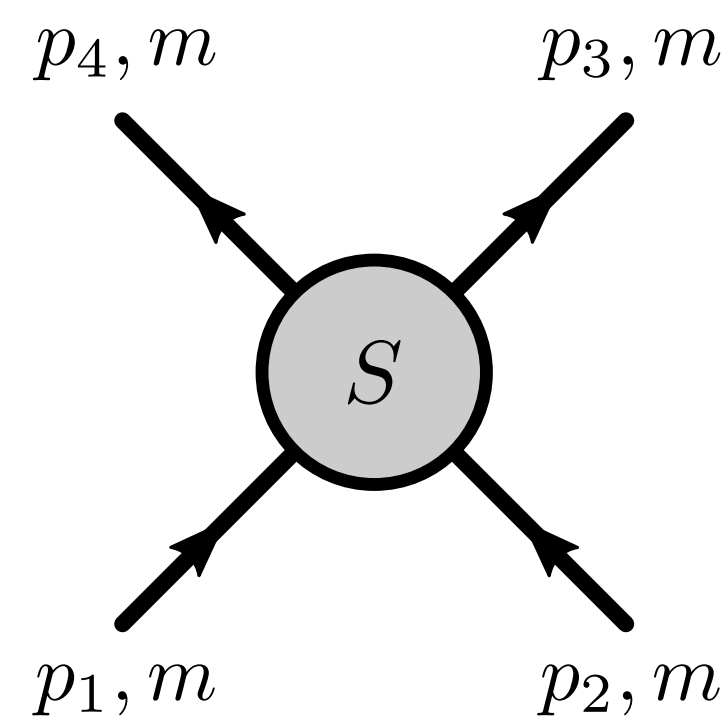
S-matrix is the evolution operator relating in and out basis of states. Conservation of probabilities

$$S^\dagger S = \mathbb{1}$$

# The standard amplitude: lightest $2 \rightarrow 2$

Most bootstrap studies focus on  $2 \rightarrow 2$  scattering of the lightest particle in the theory (*no anomalous thresholds*)

$$S = \mathbb{1} + iT \quad \langle p_3, p_4 | T | p_1, p_2 \rangle = (2\pi)^4 \delta^{(4)}(p_1 + p_2 - p_3 - p_4) T(s, t)$$



$$s = (p_1 + p_2)^2, \quad t = (p_1 - p_4)^2, \quad u = (p_1 - p_3)^2$$

$$s + t + u = 4m^2$$

## Crossing

$$T(s, t) = T(t, s) = T(u, t)$$

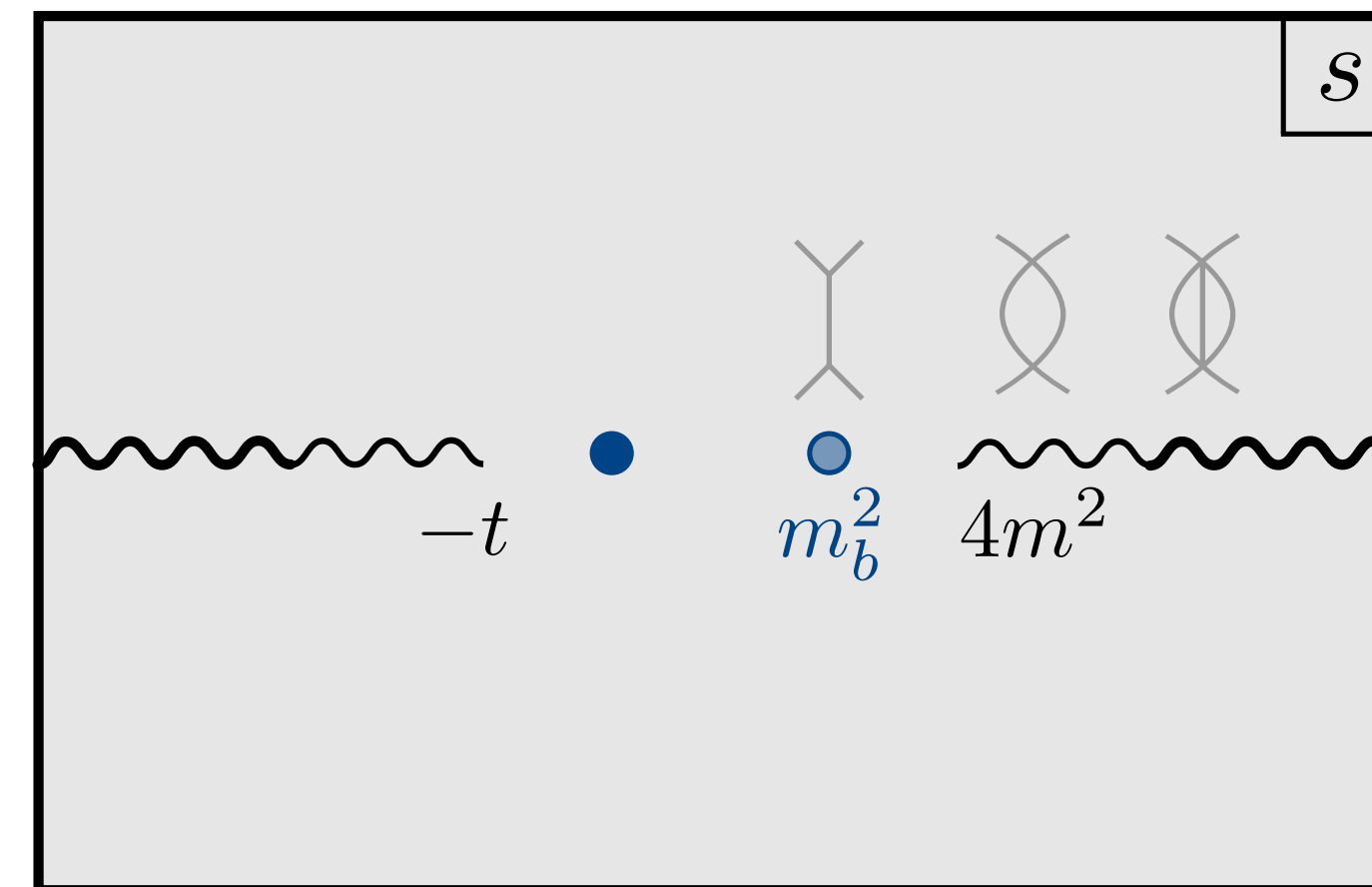
## Unitarity

$$S^\dagger S = \mathbb{1} \implies |S_{2 \rightarrow 2}|^2 + \sum_{n>2} |S_{2 \rightarrow n}|^2 = 1$$

Partial waves  $S_l(s) = 1 + i \frac{\pi}{4} \sqrt{\frac{s-4m^2}{s}} \int_{-1}^1 d \cos \theta P_l(\cos \theta) T(s, t)$

$$|S_l(s)|^2 \leq 1, \quad s > 4m^2$$

## Analyticity



## Dispersion relation (fixed t)

$$T(s, t) = \text{poles} + \int_{4m^2}^{\infty} \frac{ds'}{\pi} \frac{T_s(s', t)}{s' - s} + \int_{4m^2}^{\infty} \frac{dt'}{\pi} \frac{T_u(u', t)}{u' - u}$$

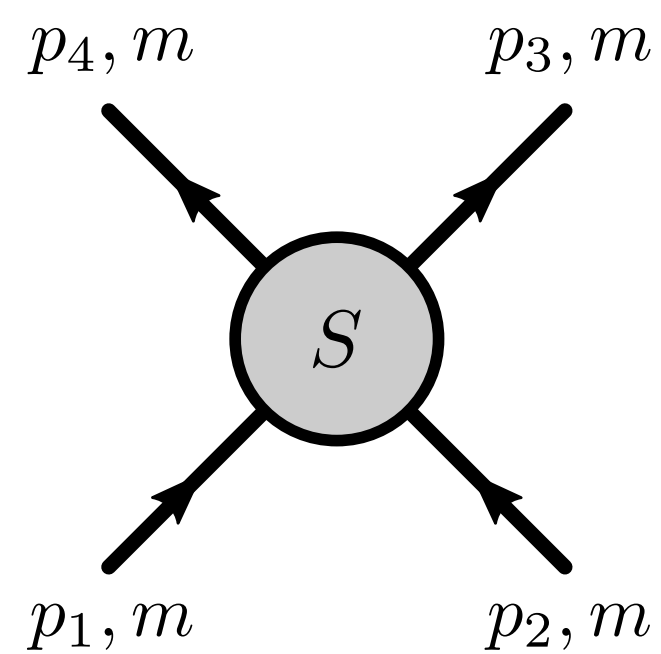
(ignoring 2 subtractions from Froissart bound  $T(s, t) \lesssim s \ln^2 s$ )

# The Methods

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# Warm-up: single particle in 2D

A simple bootstrap problem solved with complex analysis.



## Lorentz, 2D

$$S(p_1, p_2, p_3, p_4) = S(s)$$

$$s + t + u = 4m^2$$

## Crossing

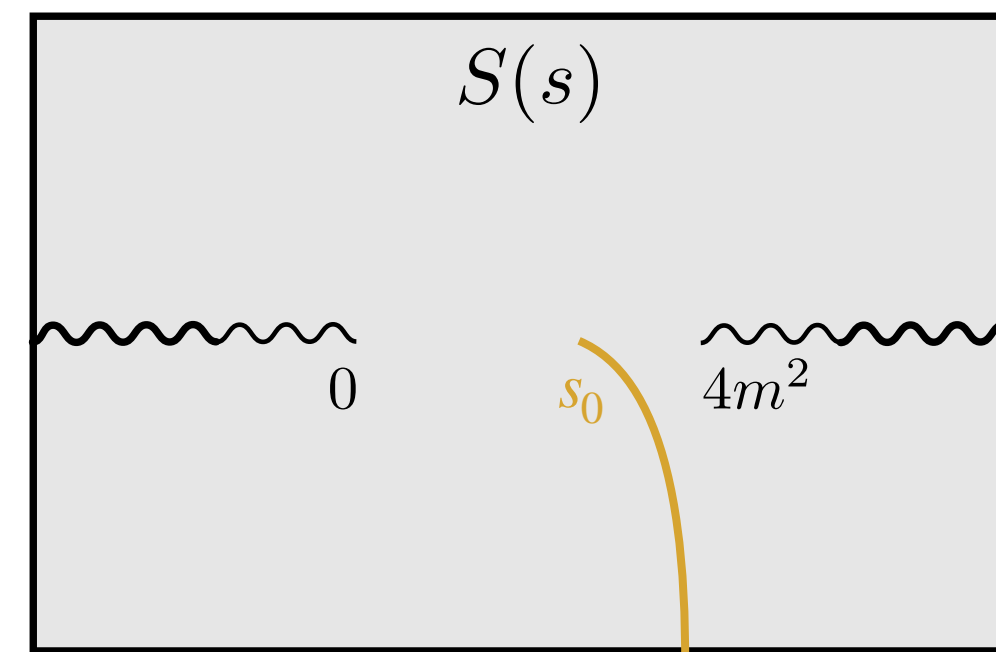
$$S(s) = S(t = 4m^2 - s)$$

## Unitarity

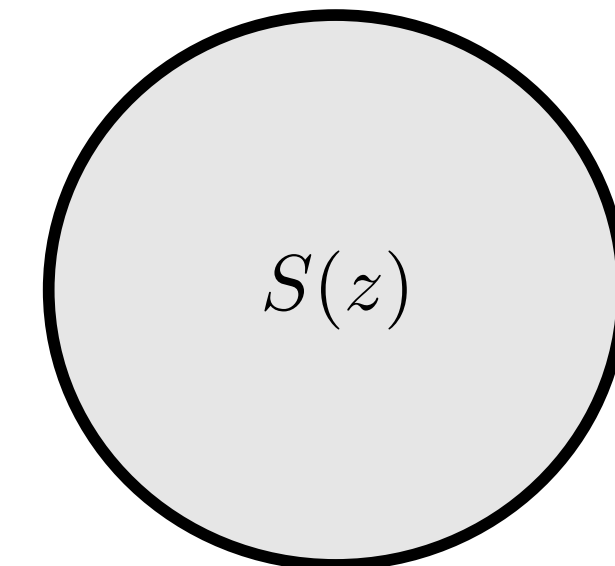
$$|S(s + i0)|^2 \leq 1, \quad s > 4m^2$$

## Analyticity

Suppose there are no bound states



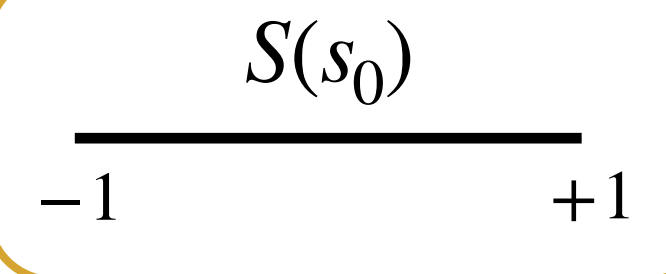
$$z(s) = \frac{\sqrt{s}\sqrt{4m^2 - s} - \sqrt{s_0}\sqrt{4m^2 - s_0}}{\sqrt{s}\sqrt{4m^2 - s} + \sqrt{s_0}\sqrt{4m^2 - s_0}}$$



What is max/min  $S(s_0)$ ?

Maximum modulus principle:  $|S(z)|$  has its maximum at the boundary (or is a constant).

→  $|S(s_0)| \leq 1$  and saturated by **free** scattering  $S = \pm 1$



# Primal approach

Constructing ACU amplitudes numerically

## Bootstrap recipe

Given a spectrum  $\{m, m_b, \dots\}$

**What is the space of allowed amplitudes?**

1. Write ansatz that trivializes **A+C**
2. Impose **U** numerically for  $s_j > 4m^2$
3. Bound parameters, i.e. max functionals  $\mathcal{F}$

## 4D analogue

1.  $T(s, t) = \sum_{a,b,c=0}^{N_{max}} \alpha_{(abc)} \rho_s^a \rho_t^b \rho_u^c$
2.  $|S_l(s)|^2 \leq 1, s_j > 4m^2, l \leq L_{max}$
3.  $\frac{(32\pi)^{-1} T(4m^2/3, 4m^2/3)}{-6.6 \quad 2.66}$

[Paulos, Penedones, Toledo, van Rees, Vieira '17]

## 2D example

$2 \rightarrow 2$ , no bound states  $\{m\}$

1. “Rho” ansatz (or dispersion relations)

$$S(s) = \sum_{n=0}^{N_{max}} \alpha_n [\rho_s^n + \rho_t^n], \quad \rho_s = \frac{\sqrt{4m^2 - s_0} - \sqrt{4m^2 - s}}{\sqrt{4m^2 - s_0} + \sqrt{4m^2 - s}}$$

2.  $|S(s_j)|^2 \leq 1, s_j > 4m^2, j \leq N_{grid}$

3. Max/min  $\mathcal{F}[S(s)] = S(2m^2)$

$$\frac{S(2m^2)}{-1 \quad +1}$$

But... bounds depend on truncation parameters  $N_{max}, L_{max} \dots$

# Dual approach

Deriving rigorous bounds from dual functionals

[2D: LC, He, Kruczenski, Vieira '19; Guerrieri, Homrich, Vieira '20; Elias, Miro, Guerrieri '21;  
4D: He, Kruczenski '21; Guerrieri, Sever '21]

$$\text{Max } \mathcal{F}[S] \leq \text{Min } \mathcal{F}_d[K]$$

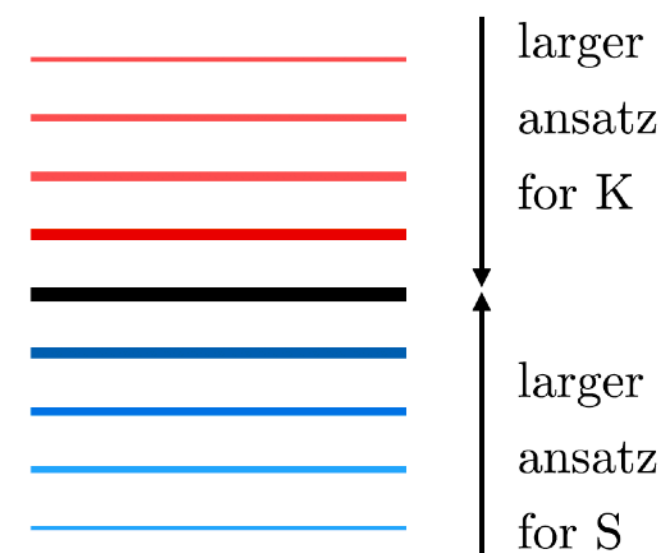
Dual functions  $K$  are Lagrange multipliers for S-matrix constraints

Simplest example:

$$\mathcal{F} = S(s_0) = \frac{1}{2\pi i} \oint_{s_0} K(s)S(s) = \frac{2}{\pi} \int_{4m^2}^{\infty} \text{Im}[K(s)S(s)] \leq \frac{2}{\pi} \int_{4m^2}^{\infty} |K(s)S(s)| \leq \frac{2}{\pi} \int_{4m^2}^{\infty} |K(s)| = \mathcal{F}_d$$

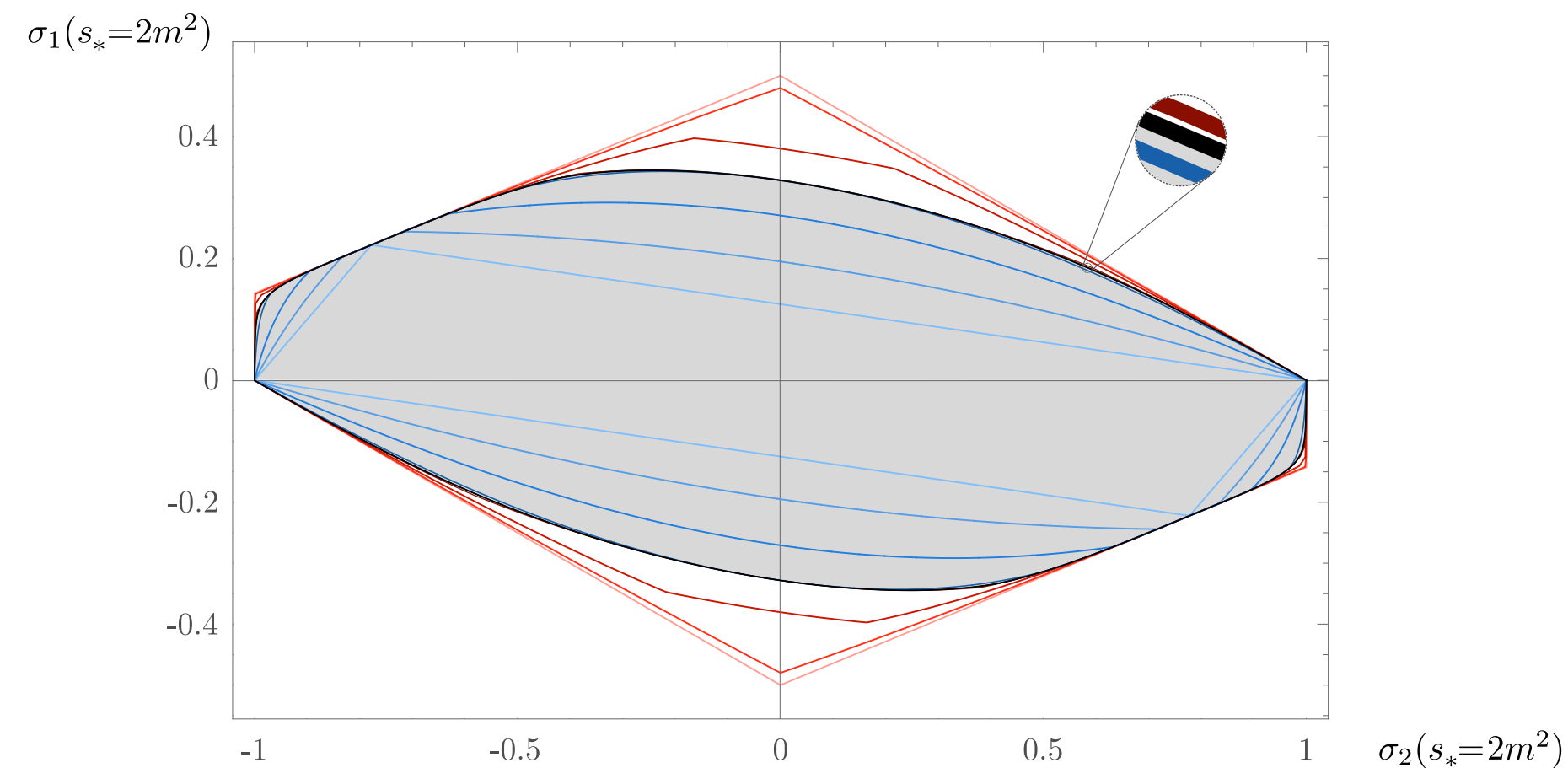
$K(s) \sim \frac{1}{s-s_0}$        $K(4m^2-s) = -K(s)$        $|S(s)| \leq 1$

Run Bootstrap for  $\mathcal{F}_d[K]$

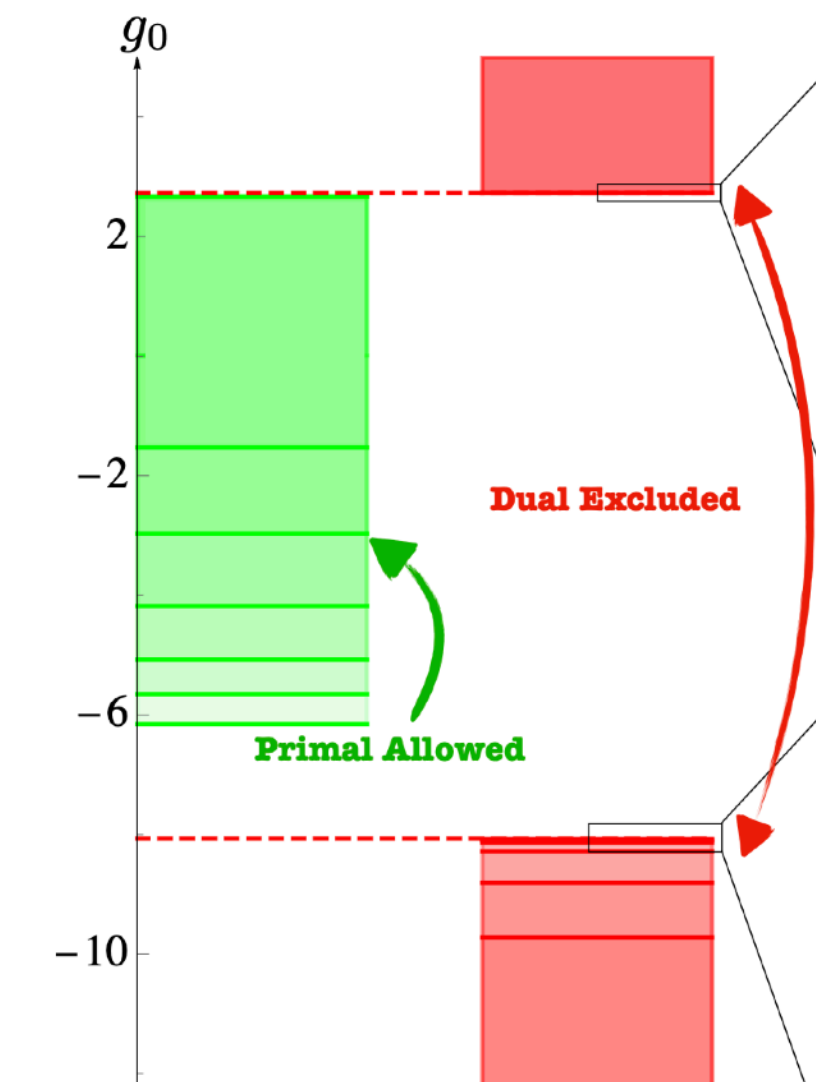


optimality

$$S(s) = \frac{iK^*(s)}{|K(s)|}$$



2D,  $O(N)$  sym [LC, He, Kruczenski, Vieira '19]



4D,  $T(4m^2/3, 4m^2/3)$   
[Guerrieri, Sever '21]

# Applications

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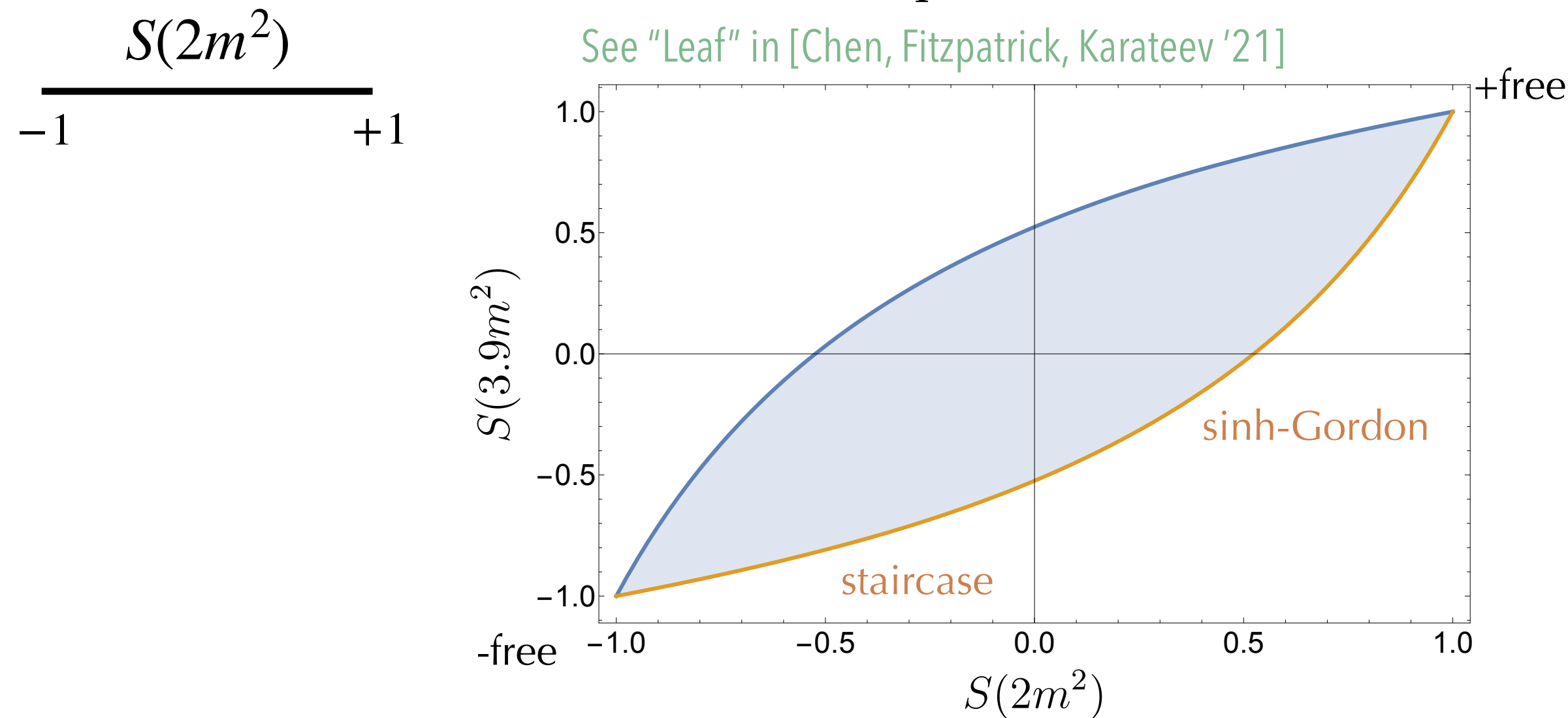
# Maps in 2D

Two-dimensional scattering is much simpler, but still has a lot to teach us!

## Revisiting warm-up

What if we explore more directions?

See "Leaf" in [Chen, Fitzpatrick, Karateev '21]



Bounds saturated by CDD zeros: resonances

$$S(s) = \pm \frac{\sqrt{s}\sqrt{4m^2 - s} - \sqrt{s_0}\sqrt{4m^2 - s_0}}{\sqrt{s}\sqrt{4m^2 - s} + \sqrt{s_0}\sqrt{4m^2 - s_0}}$$

$$S(s) = S(4m^2 - s) \quad |S(s)|^2 = 1$$

No particle production  $\rightarrow$  integrable models

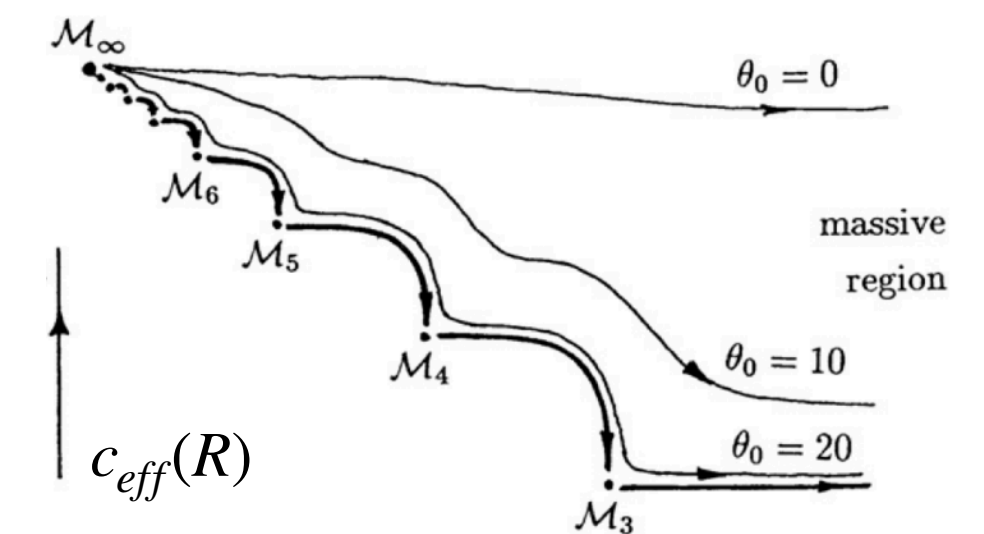
sinh-Gordon

$$0 \leq s_0 \leq 4m^2$$

staircase

$$s_0 = 2m^2 + i\mathbb{R}$$

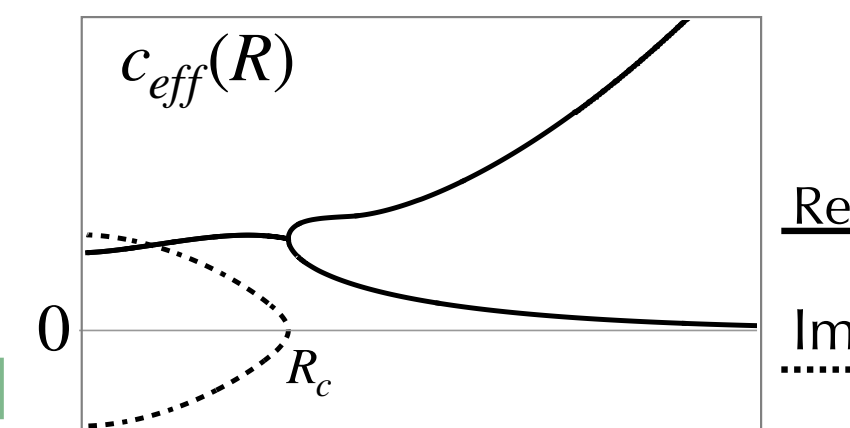
$$\mathcal{L}_{int}^{shG} \supset \frac{m_0^2}{\beta^2} \cosh \beta \phi$$



From TBA: [Zamolodchikov '06]

Add more directions  $\rightarrow \prod$  CDD zeros

“Consistent” amplitudes with many resonances  
 $\rightarrow$  unconventional UV behaviour

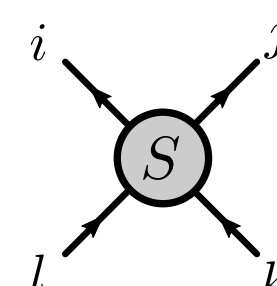


[Camilo, Fleury, Lencses, Negro, Zamolodchikov '21; LC, Negro, Schaposnik Massolo '21]

# 2D: Global group symmetries

More analytic functions, integrability non-generic

## O(N) monolith



$$S_{lk}^{ij}(s) = \sigma_1(s) \begin{array}{c} i \quad j \\ \text{---} \\ l \quad k \end{array} + \sigma_2(s) \begin{array}{c} i \quad j \\ \text{---} \\ l \quad k \end{array} + \sigma_3(s) \begin{array}{c} i \quad j \\ \text{---} \\ l \quad k \end{array}$$

$$= S_{\bullet}(s) (\mathbb{P}_{\bullet})_{ij}^{kl} + S_A(s) (\mathbb{P}_A)_{ij}^{kl} + S_S(s) (\mathbb{P}_S)_{ij}^{kl}$$

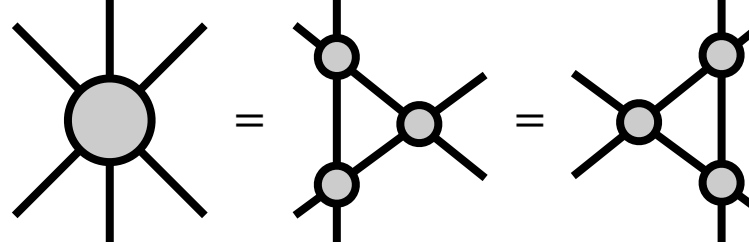
**Crossing**

$$S_a(4m^2 - s) = \sum_{a'} C_{a,a'} S_{a'}(s)$$

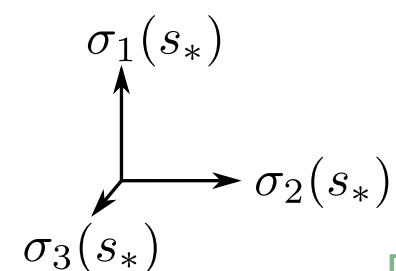
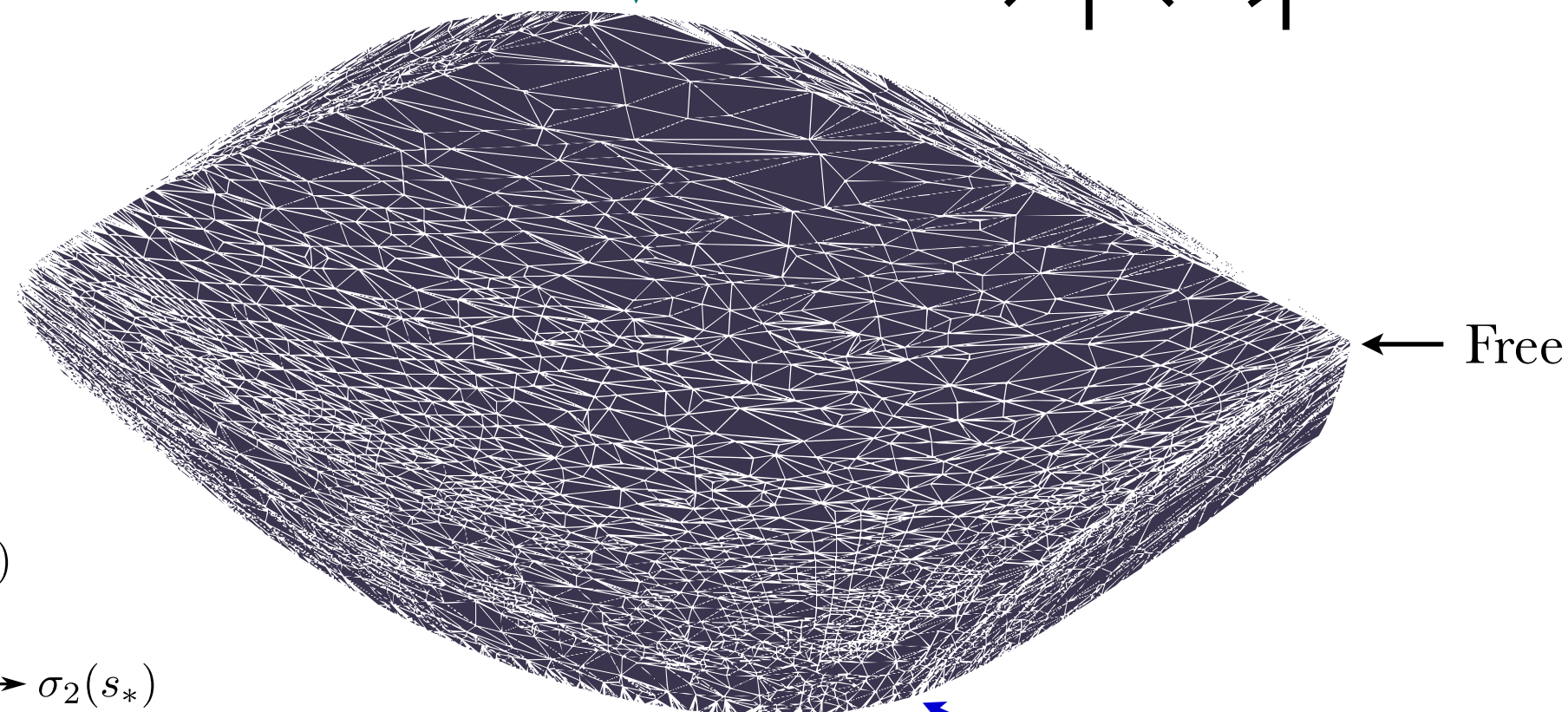
**Unitarity**

$$|S_a(s^+)| \leq 1$$

Yang-Baxter equations



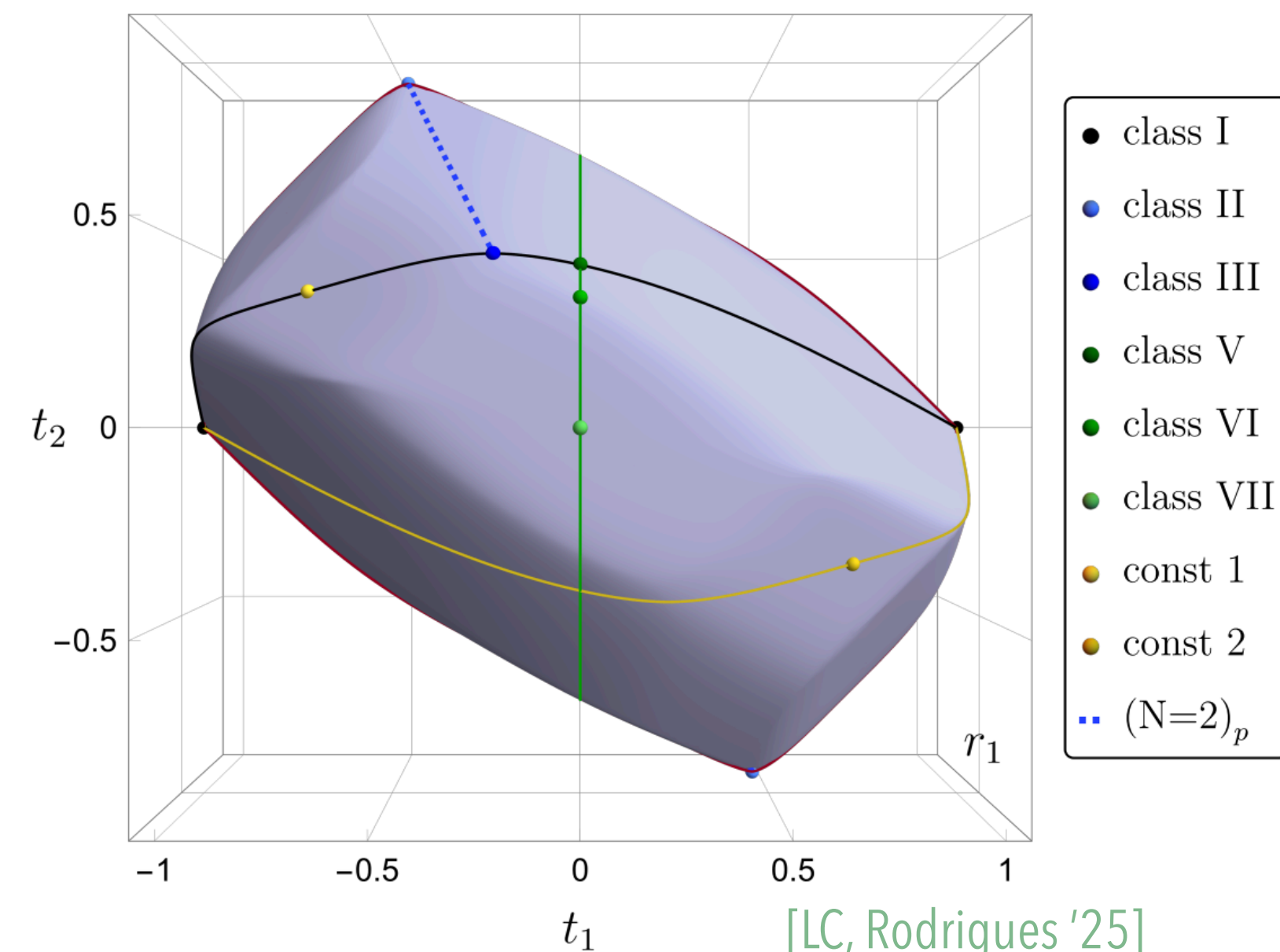
periodic YB



[LC, He, Kruczenski, Vieira '19]

## U(N) monolith

6 functions of  $s$ :  $t_{1,2}$ ,  $r_{1,2}$ ,  $u_{1,2}$



[LC, Rodrigues '25]

- Generic: unitarity saturation (even w/o YB), emergent periodicity in  $\theta \sim \ln s$
- Integrable periodic amplitudes show 'walking' behaviour, possible connection to complex CFTs [Gorbenko, Zan '20]

# Generalized Symmetries

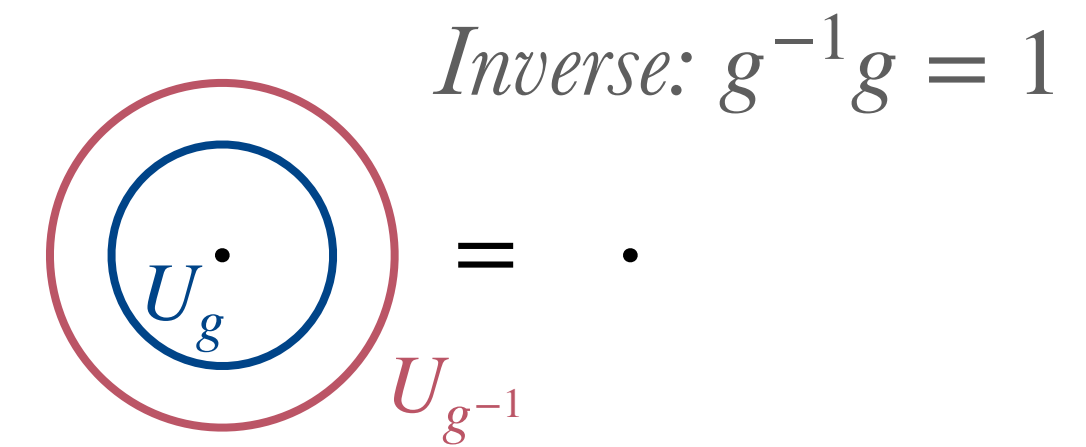
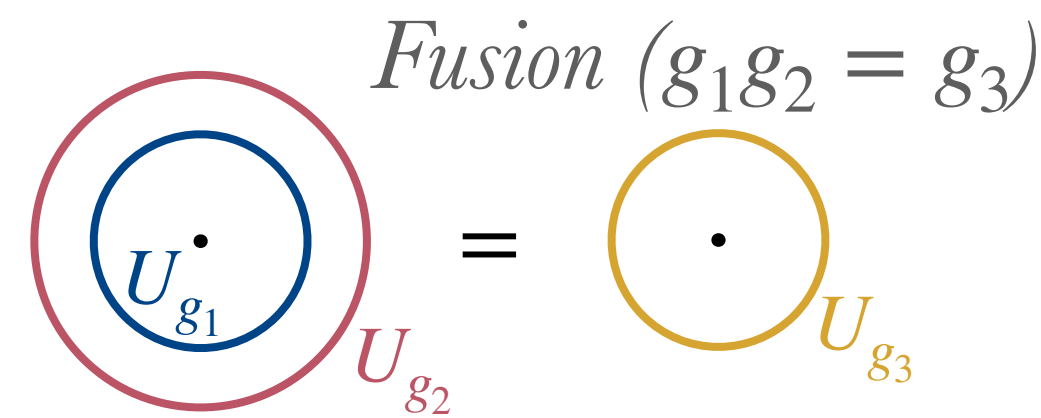
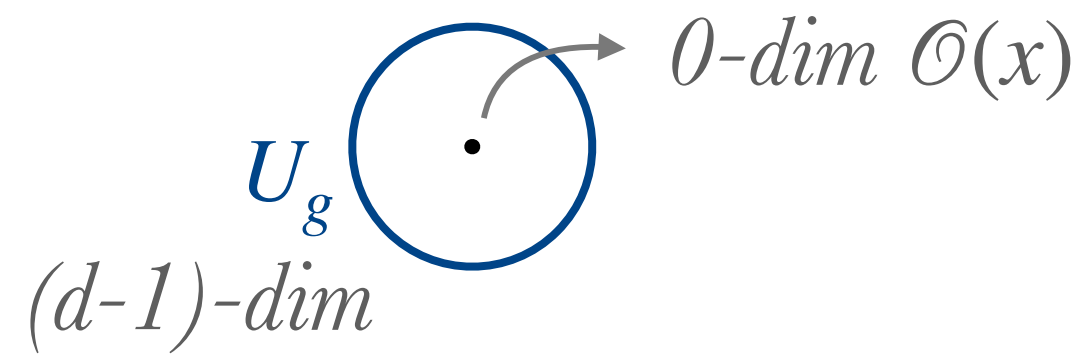
- (global) Symmetries in QFT  $\leftrightarrow$  Topological operators [Gaiotto, Kapustin, Seiberg, Willett '14]

*Move in  $t \rightarrow$  charge conservation*

*e.g. from Noether current  $j_\mu$ :  $U = \exp\left(ia \oint d^{d-1}x j_0(x)\right)$*



- Usual group symmetry: 0-form, invertible

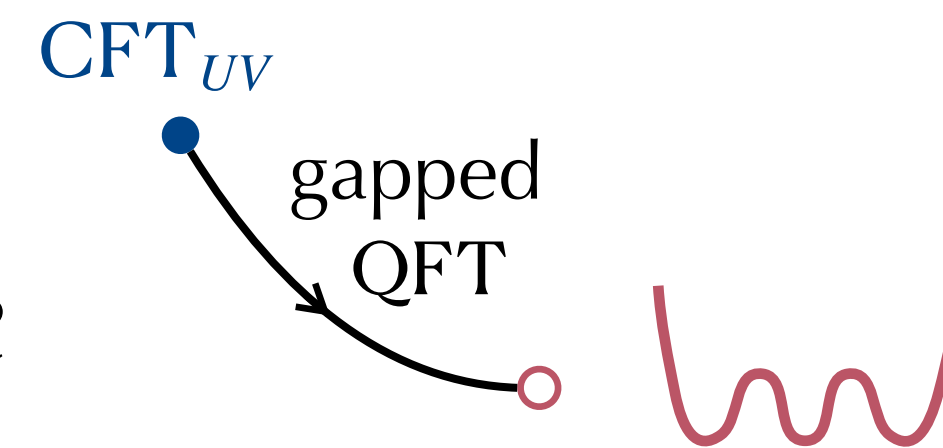


- Generalizations: Higher-form symmetries.  $q$ -form sym,  $(d-q-1)$ -dim topological ops.  
Non-invertible symmetries.  $g^{-1}$  not necessary

Here: 0-form symmetries in 2d  $\rightarrow$  topological lines  $\mathcal{L}_a$  described by **Fusion categories** (e.g. Ising)

- $\exists$  symmetries and anomalies  $\rightarrow$  strong constraints on IR

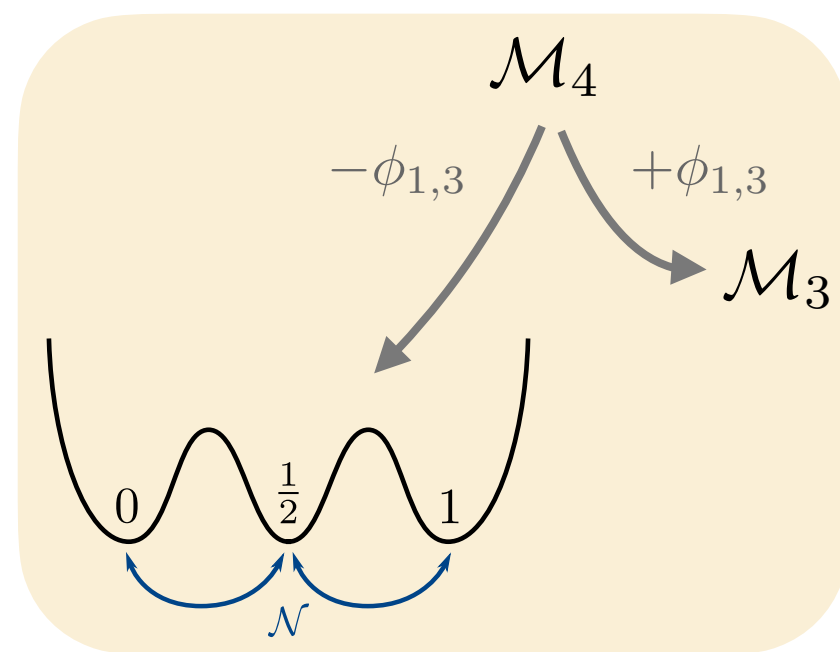
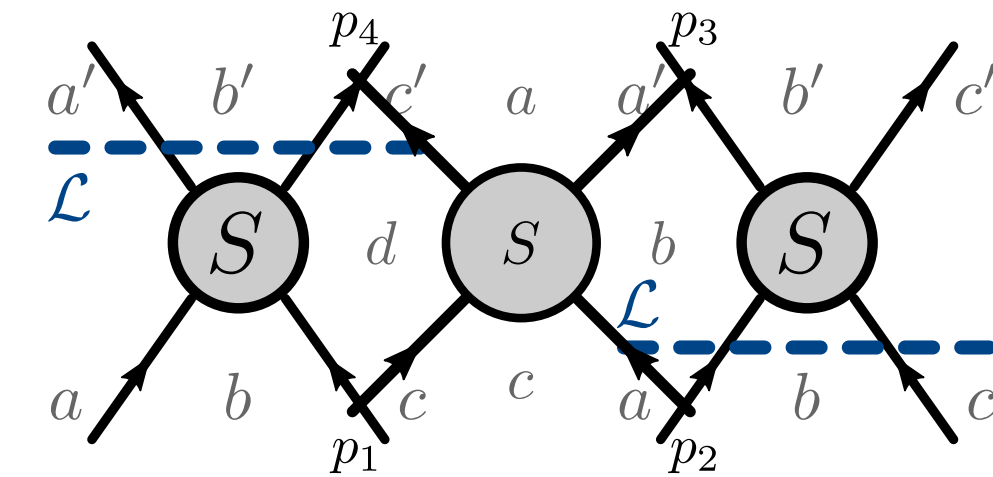
e.g. **degenerate vacua** in IR  
[Chang, Lin, Shao, Wang, Yin '18]



# Non-invertible symmetries and modified crossing

Scattering of kinks interpolating between vacua

Implications of symmetries: **Crossing symmetry of S-matrix gets modified!**



Example: integrable deformation of Tricritical Ising 1+1d

- unitarity + **crossing** + **integrability** [Zamolodchikov '89] → ~~non-invertible syms~~

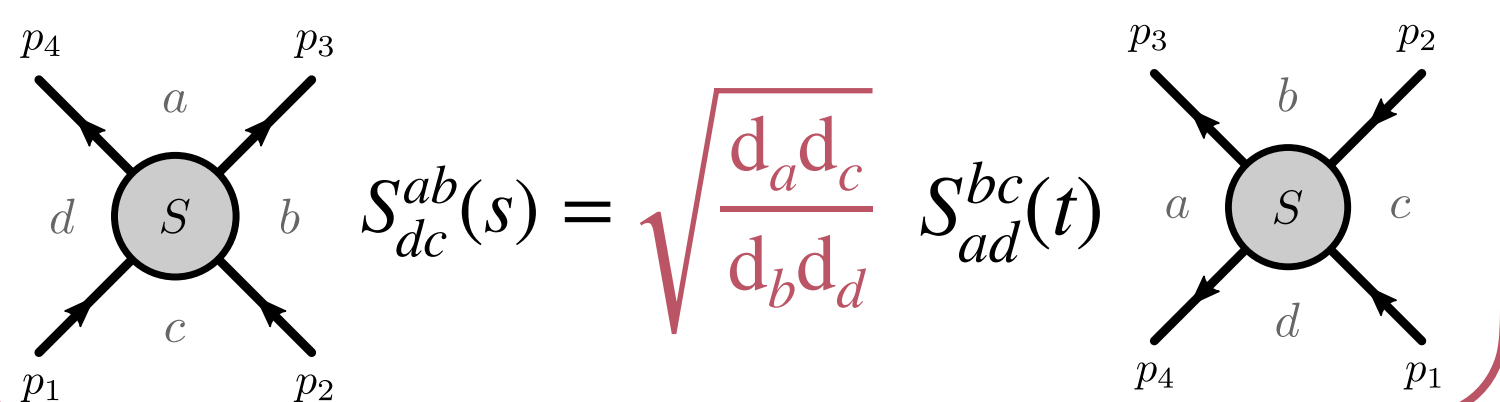
$$\widehat{S}_{dc}^{ab}(\theta) = Z(\theta) \left( \frac{d_a d_c}{d_b d_d} \right)^{\frac{i\theta}{2\pi}} \left[ \sqrt{\frac{d_a d_c}{d_b d_d}} \sinh\left(\frac{\theta}{n}\right) \delta_{bd} + \sinh\left(\frac{i\pi - \theta}{n}\right) \delta_{ac} \right]$$

$(s = 4m^2 \cosh^2(\theta/2), n = 4, d_0 = d_1 = 1, d_{1/2} = \sqrt{2})$

- unitarity + **non-invertible syms** + **integrability**

$$S_{dc}^{ab}(\theta) = Z(\theta) \left[ \sqrt{\frac{d_a d_c}{d_b d_d}} \sinh\left(\frac{\theta}{n}\right) \delta_{bd} + \sinh\left(\frac{i\pi - \theta}{n}\right) \delta_{ac} \right]$$

→ **Modified crossing** [Copetti, LC, Komatsu '24]



symmetry data ( $d_a$ : quantum dimensions)

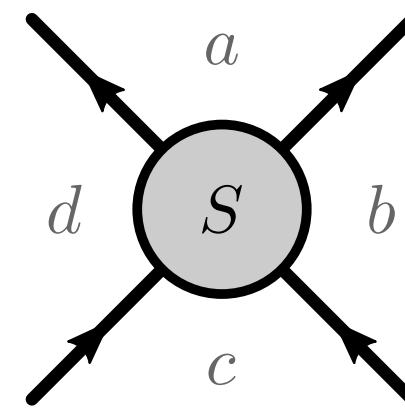
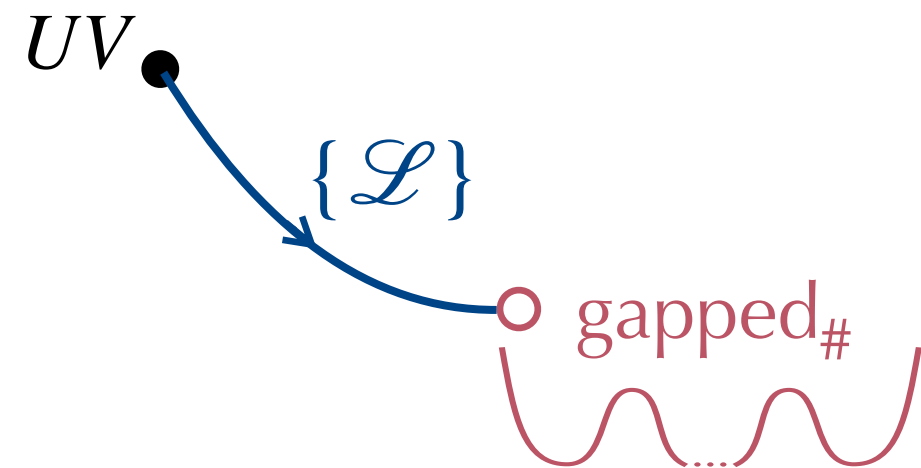
In general: normalization of  $|in\rangle$  and  $|out\rangle$  states, taking into account topological degrees of freedom.

$$\langle in|in\rangle_s = d \begin{matrix} \text{---} c \text{---} \\ | \quad | \\ v \quad v \\ | \quad | \\ \text{---} c \text{---} \end{matrix} b \times \delta^2(\cdot)$$

# S-matrix Bootstrap with non-inv. sym.

**Set-up:**

2D QFT with categorical sym.



2 → 2 Kink scattering

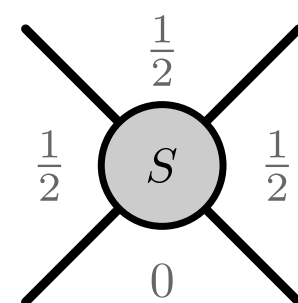
Bootstrap from **A+mC+U+S** assuming:

- { minimum spectrum required by symmetry
- { SSB all symmetries (#vacua=#L)

► **Symmetries**  $\{\mathcal{L}\}$   $S_{dc}^{ab}(s) = \sum_{\chi} A_{\chi}(s)(P_{\chi})_{dc}^{ab} \rightarrow$  Projector into fusion channel  $\chi$

e.g. **Fibonacci category** ( $W^2 = 1+W$ )

- 2 vacua  $a = 0, 1/2$
- 2 partial amplitudes  $A_0(s), A_{1/2}(s)$



$K_{a,a\pm 1/2}$   
 $K_{1/2,1/2}$

• Cubic coupling

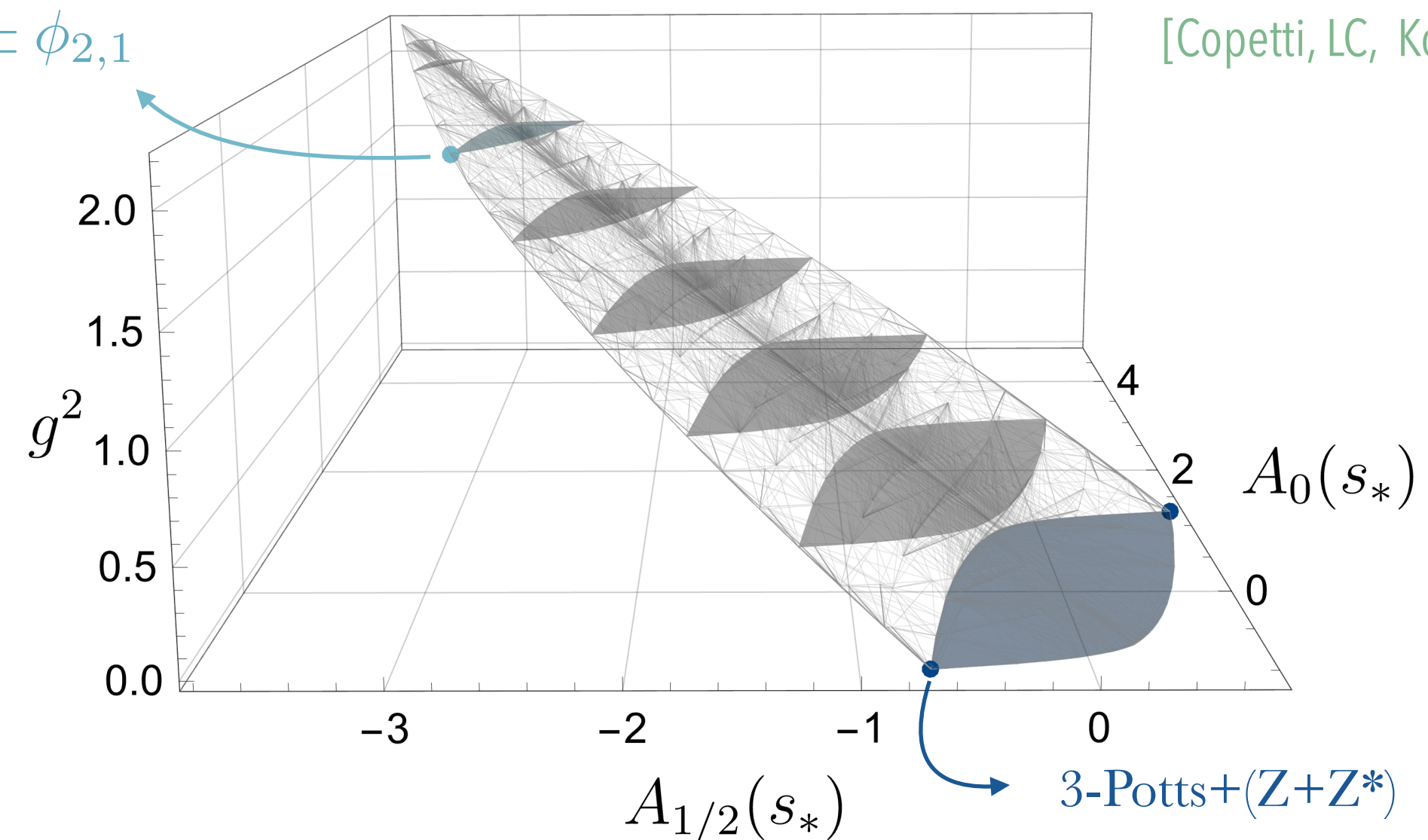
$$A_{1/2}(s) \sim \frac{g^2}{s - m^2}$$

Bootstrap  $\{A_0(s_*), A_{1/2}(s_*), g \sim \text{diagram}\} \Rightarrow$

Fibonacci

$\mathcal{M}_4 \pm \phi_{2,1}$

[Copetti, LC, Komatsu '24]



# 4D: Pions

Pions  $\pi^{\pm,0}$  pseudo-Goldstone bosons for chiral symmetry breaking, triplet  $SU(2)_V$ .

$$\pi_a(p_1) + \pi_b(p_2) \rightarrow \pi_c(p_3) + \pi_d(p_4) \quad T_{ab,cd} = A(s, t, u)\delta_{ab}\delta_{cd} + A(t, s, u)\delta_{ac}\delta_{bd} + A(u, t, s)\delta_{ad}\delta_{bc}$$

Two recent approaches:

## Adding UV info through Form Factors

[He, Kruczenski '23-25]

Include currents  $\int e^{ipx} \mathcal{O}(x) |0\rangle$

[Karateev, Kuhn, Penedones '19]

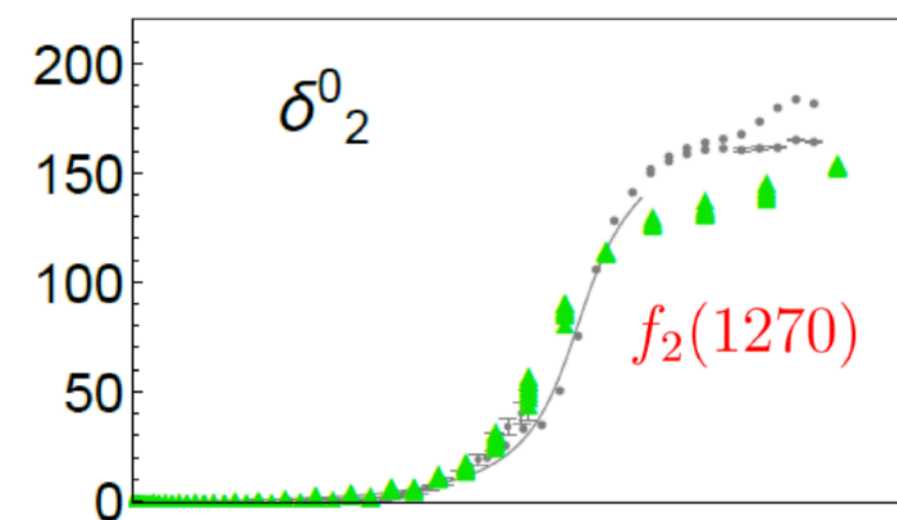
$$\begin{pmatrix} 1 & S^* & F^* \\ S & 1 & F \\ F & F^* & 2\pi\rho \end{pmatrix} \geq 0$$

Energy

Pert. QCD  
FF asymptotics

S-matrix/FF  
Bootstrap

Chiral SB



$\Rightarrow$  Reproduce masses for  $\rho, \rho', f_2$

## Fitting to experimental data

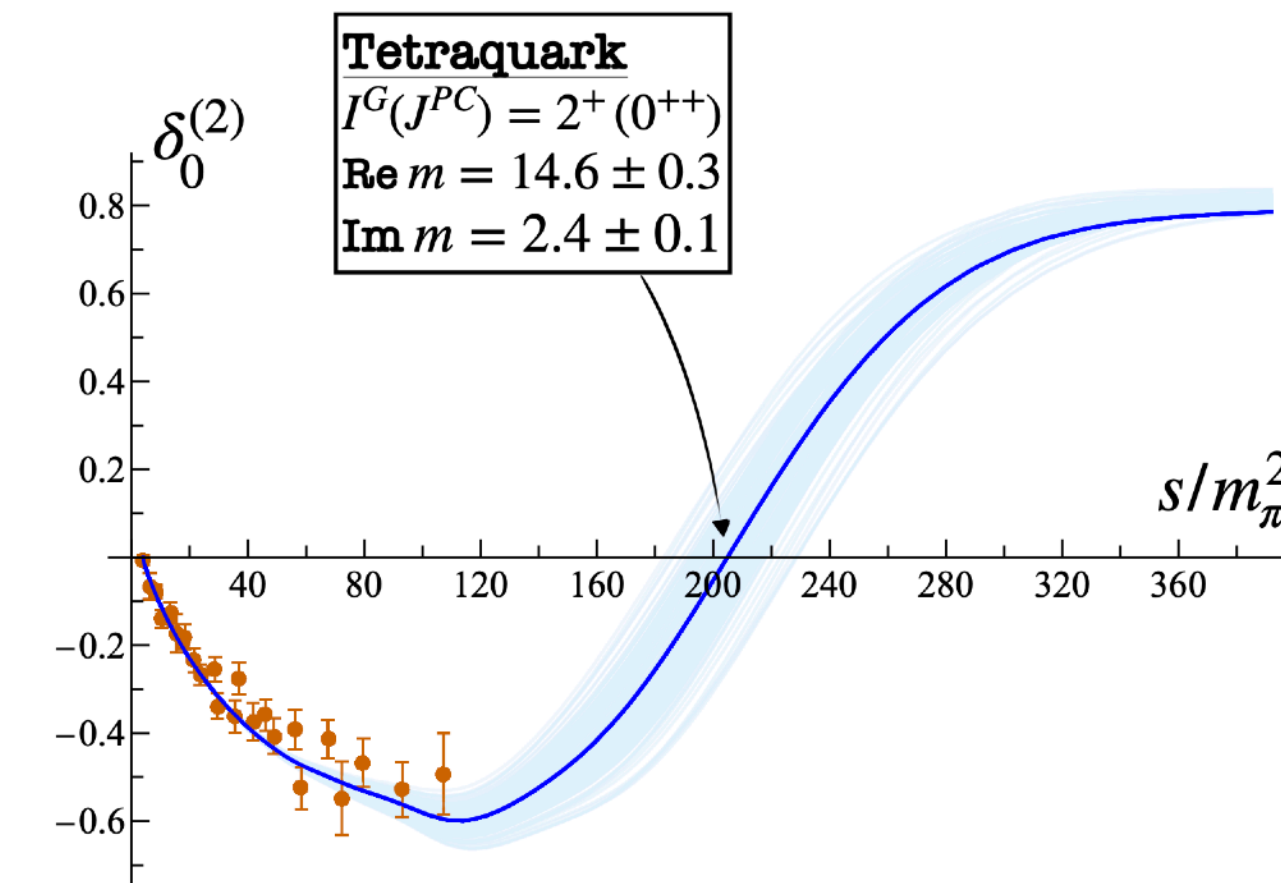
[Guerrieri, Häring, Su '24]

Input e.g. masses for  $\rho, f_0, f_0', f_2$

Energy

S-matrix  
Bootstrap

Fit to data



$\Rightarrow$  Predict new resonances

# Final Remarks

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# Final Remarks

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## S-matrix Bootstrap

A very useful tool to carve out the space of consistent QFTs/learn more about your favourite theory

Many applications I did not talk about: e.g. Effective Field Theories and Gravity (see rest of the conference)

Some future directions I am interested in :

- Beyond lightest  $2 \rightarrow 2$  scattering Multiparticle Flux-Tube [Guerrieri, Homrich, Vieira '24]  
Non-perturbative anomalous thresholds [Correia '22]
- Axioms not fully proven, might need modifications
- 4D gauge and gravity: IR finite observables for long-range interactions?

# Thank you!