

Crash course on tools for Dark Matter research

Roberto A. Lineros

Space sciences, Technologies and Astrophysics Research (STAR) Institute
Université de Liège

Dark Matter Days 2017 – CIFFU BUAP

Motivation

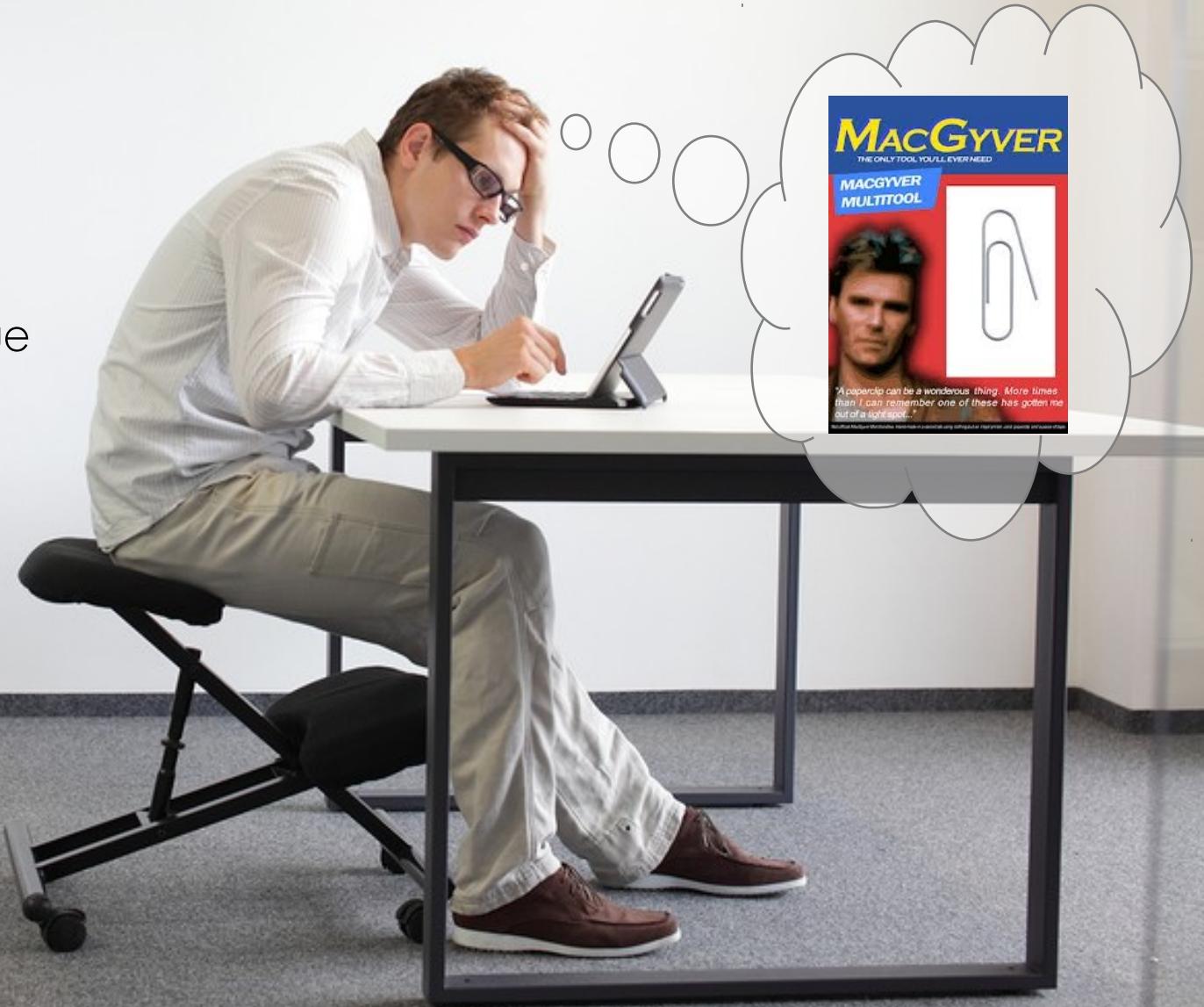
Research on Dark Matter
could be hard without
a good set tools



disclaimer: I don't know who is this guy

Motivation

We hope for a unique
tool to solve all our
problems (in physics)

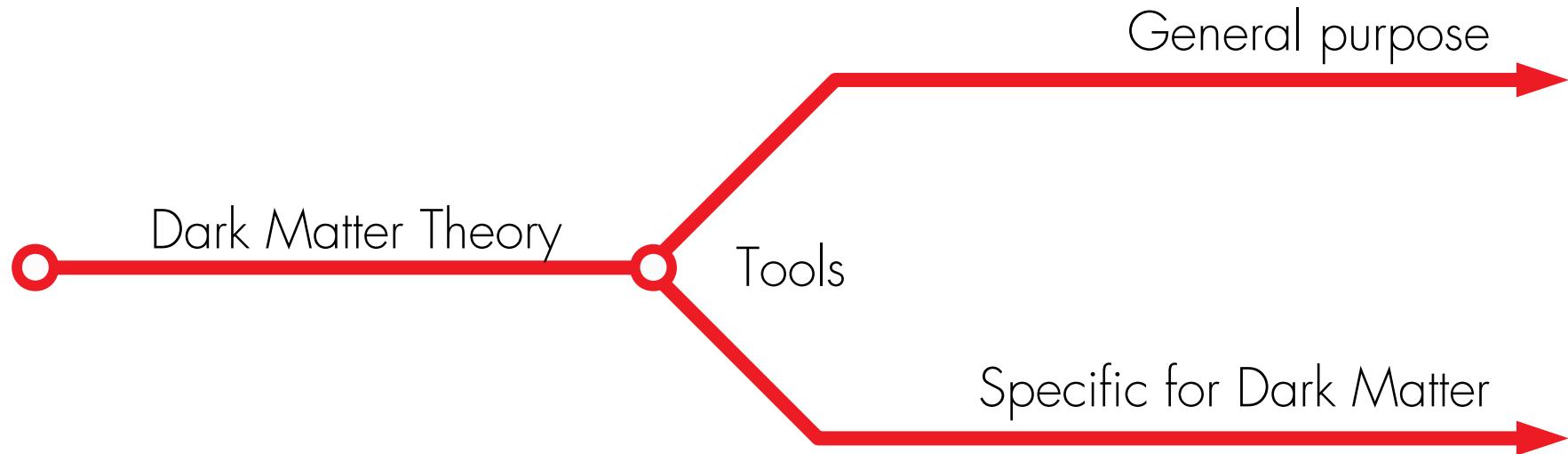


Motivation

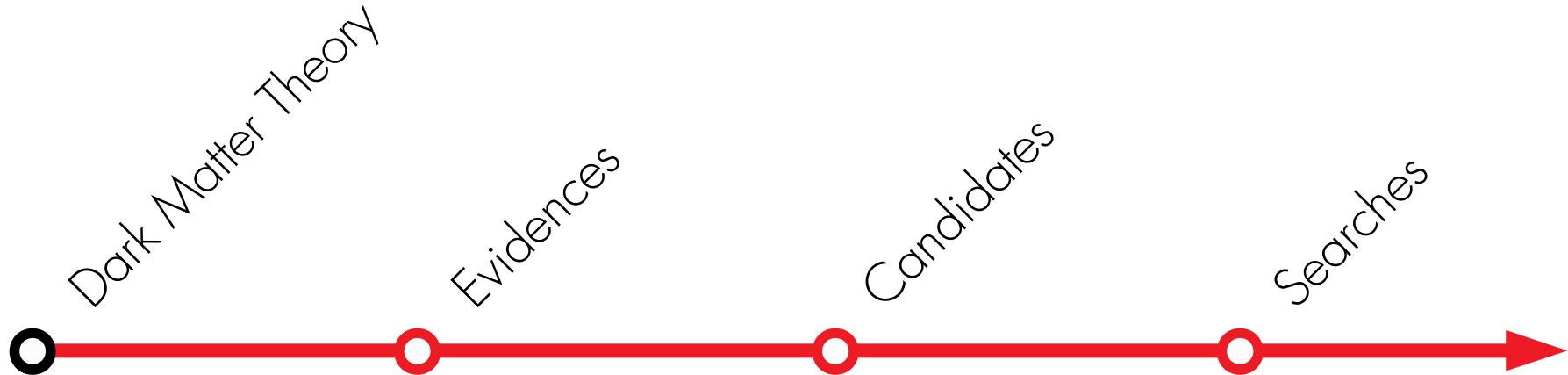
... but most of the tools are made by physicists to solve their own problems



Course plan



Today's plan





Dark Matter: Evidences

A little of history



1933: **Fritz Zwicky** postulated dark matter in order to explain **Coma cluster dynamics** as a bound system



<http://youtu.be/C8ZOgKlsqEA>

A little of history

THE ASTROPHYSICAL JOURNAL

AN INTERNATIONAL REVIEW OF SPECTROSCOPY AND
ASTRONOMICAL PHYSICS

VOLUME 86

OCTOBER 1937

NUMBER 3

ON THE MASSES OF NEBULAE AND OF CLUSTERS OF NEBULAE

F. ZWICKY

ABSTRACT

Present estimates of the masses of nebulae are based on observations of the *luminosities* and *internal rotations* of nebulae. It is shown that both these methods are unreliable; that from the observed luminosities of extragalactic systems only lower limits for the values of their masses can be obtained (sec. i), and that from internal rotations alone no determination of the masses of nebulae is possible (sec. ii). The observed internal motions of nebulae can be understood on the basis of a simple mechanical model, some properties of which are discussed. The essential feature is a central core whose internal *viscosity* due to the gravitational interactions of its component masses is so high as to cause it to rotate like a solid body.

In sections iii, iv, and v three new methods for the determination of nebular masses are discussed, each of which makes use of a different fundamental principle of physics.

Method iii is based on the *virial theorem* of classical mechanics. The application of this theorem to the Coma cluster leads to a minimum value $\bar{M} = 4.5 \times 10^{10} M_{\odot}$ for the average mass of its member nebulae.

Method iv calls for the observation among nebulae of certain *gravitational lens* effects.

Section v gives a generalization of the principles of ordinary *statistical mechanics* to the whole system of nebulae, which suggests a new and powerful method which ultimately should enable us to determine the masses of all types of nebulae. This method is very flexible and is capable of many modes of application. It is proposed, in particular, to investigate the distribution of nebulae in individual great clusters.

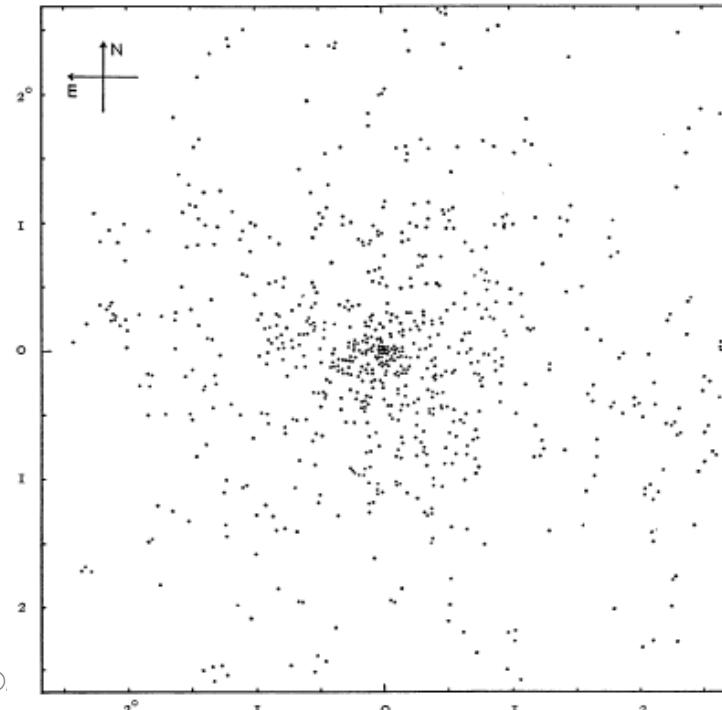
As a first step toward the realization of the proposed program, the Coma cluster of

THE MASSES OF NEBULAE

227

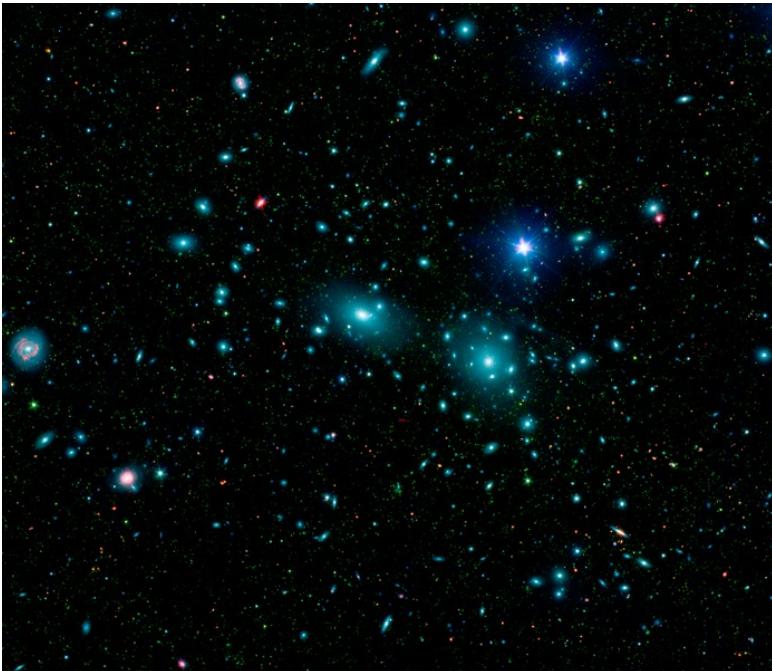
III. THE VIRIAL THEOREM APPLIED TO CLUSTERS OF NEBULAE

If the total masses of clusters of nebulae were known, the average masses of cluster nebulae could immediately be determined from counts of nebulae in these clusters, provided internebular material is of the same density inside and outside of clusters.

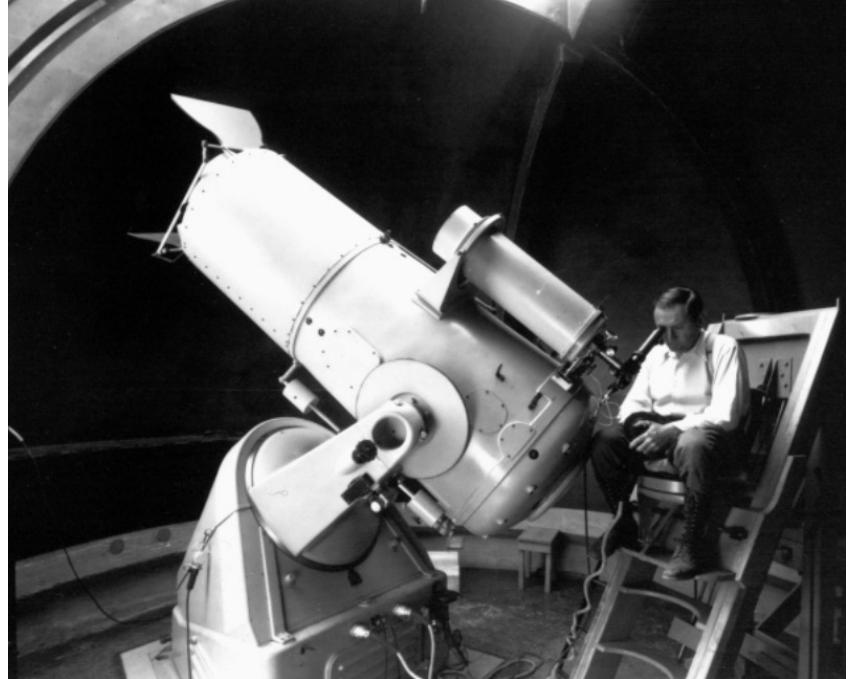


Crash course on D.

A little of history



Coma Galaxy Cluster



Fritz Zwicky

Virial Theorem



Dark Matter

Virial theorem

It relates global quantities
of a system in equilibrium

$$\langle T \rangle = -\frac{1}{2} \sum_{k=1}^N \langle \vec{F}_k \cdot \vec{r}_k \rangle$$

For gravitationally bounded objects:

$$\langle T \rangle = -\frac{1}{2} \langle V_{\text{total}} \rangle \quad \longrightarrow \quad M_{\text{total}} \simeq \frac{2R_{\text{total}} \bar{v}^2}{G}$$

A little of history



60's – 70's **Vera Rubin** studied rotation curves of many galaxies



Andromeda galaxy

Doppler effect and redshift

Observed redshift



$$z = \frac{\lambda_{\text{obs}} - \lambda_{\text{emit}}}{\lambda_{\text{emit}}}$$

$$1 + z = \frac{1 + v \cos \theta / c}{\sqrt{1 + v^2 / c^2}}$$

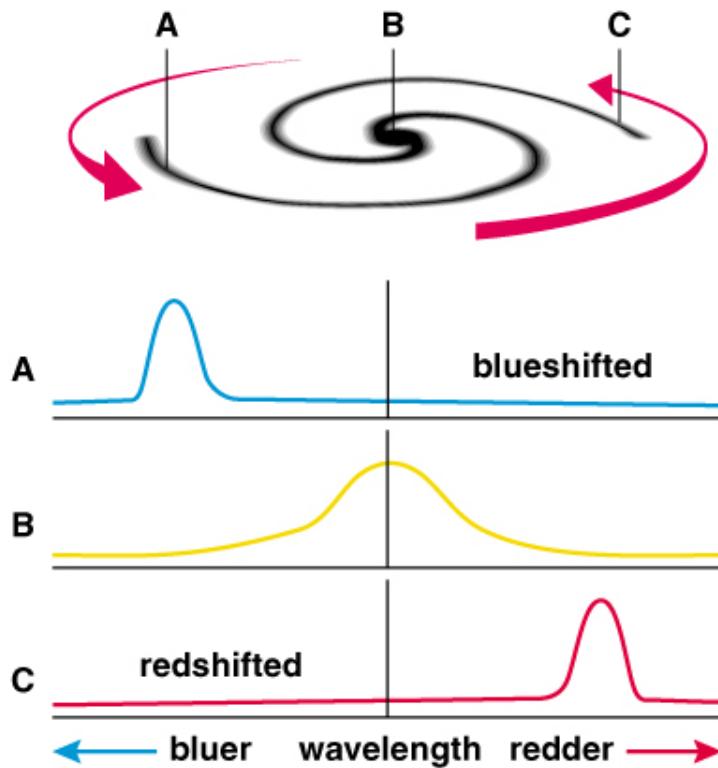


Estimation of the velocity

$z > 0$ object going away

$z < 0$ object getting closer

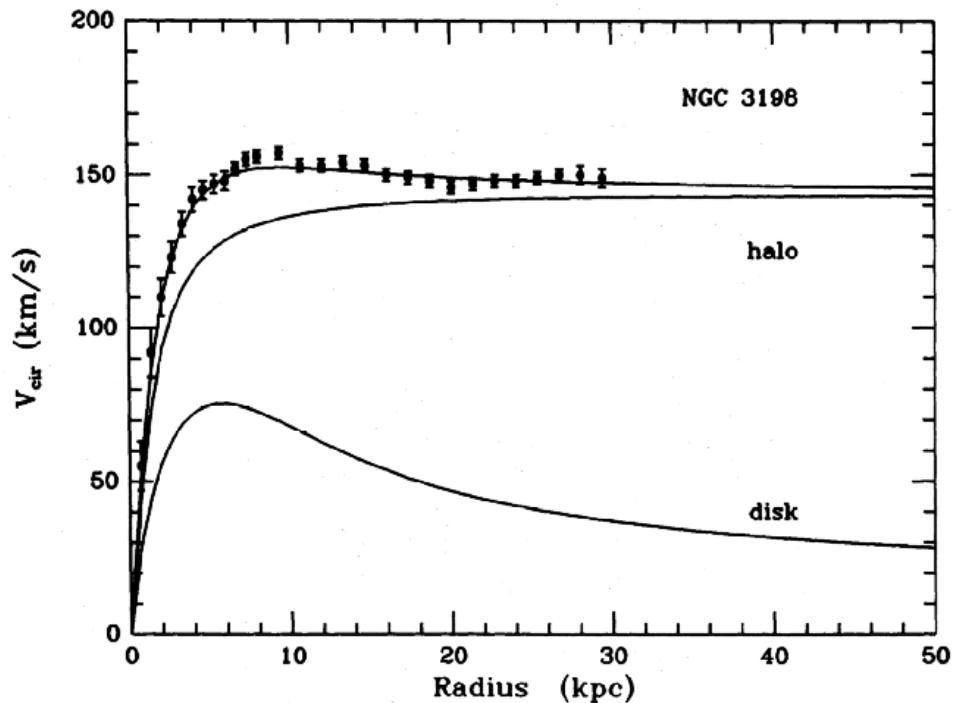
A little of history



Copyright © Addison Wesley.

8 Nov 2017

60's – 70's Vera Rubin studied rotation curves of many galaxies



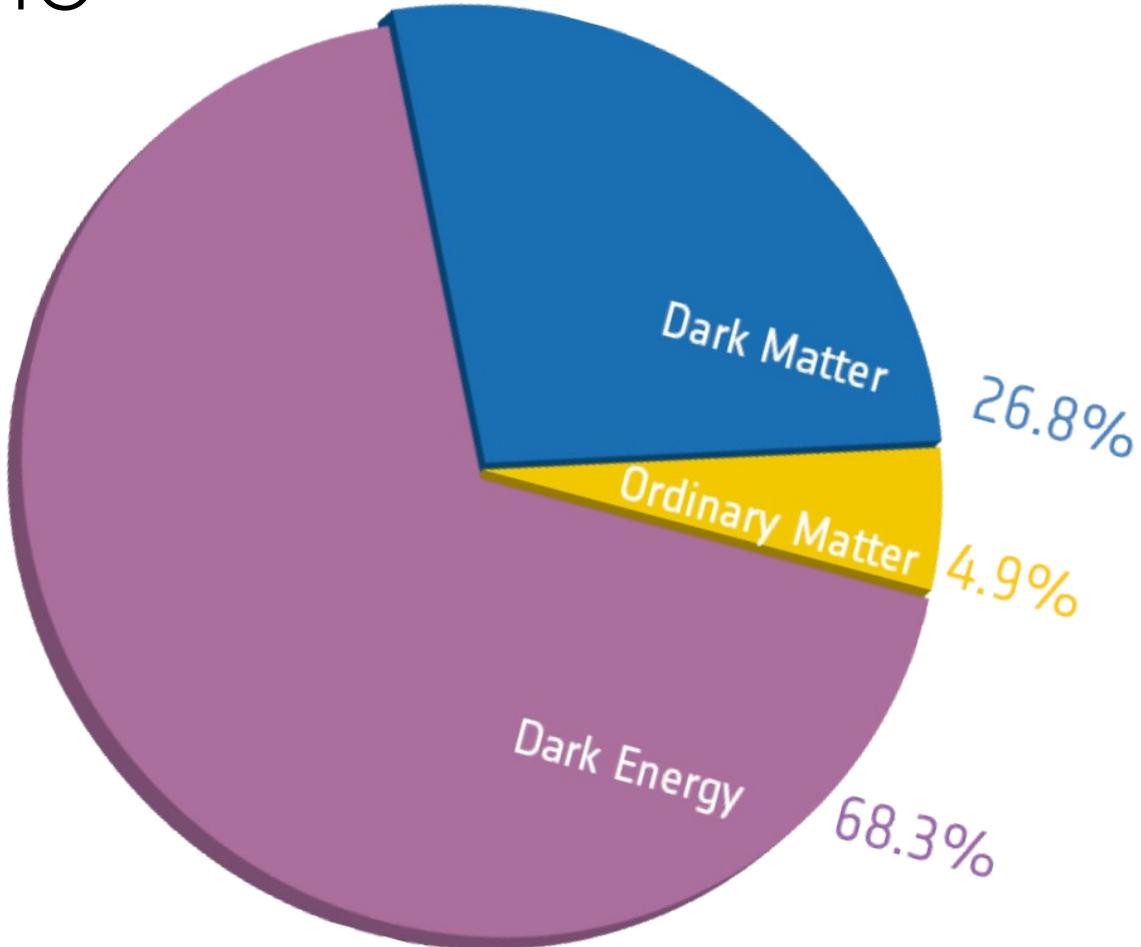
R. A. Lineros. Crash course on DM

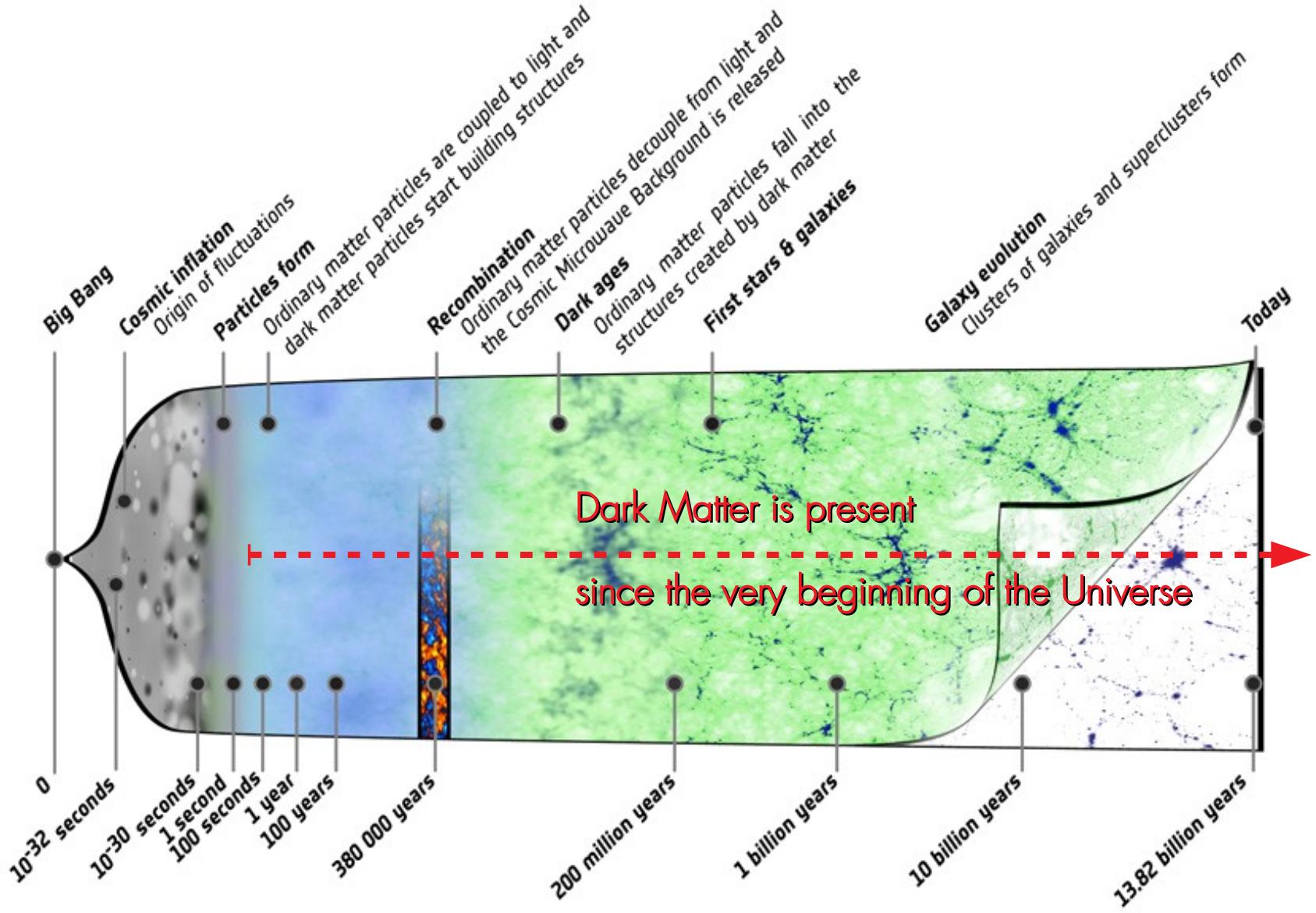
14/62

The background of the image is a deep space scene filled with numerous galaxies of various sizes and colors, primarily in shades of yellow, orange, and white. A prominent, larger yellow galaxy is visible in the upper left corner. In the center-right area, there is a bright, multi-pointed starburst or lens flare effect. The overall atmosphere is dark and star-filled.

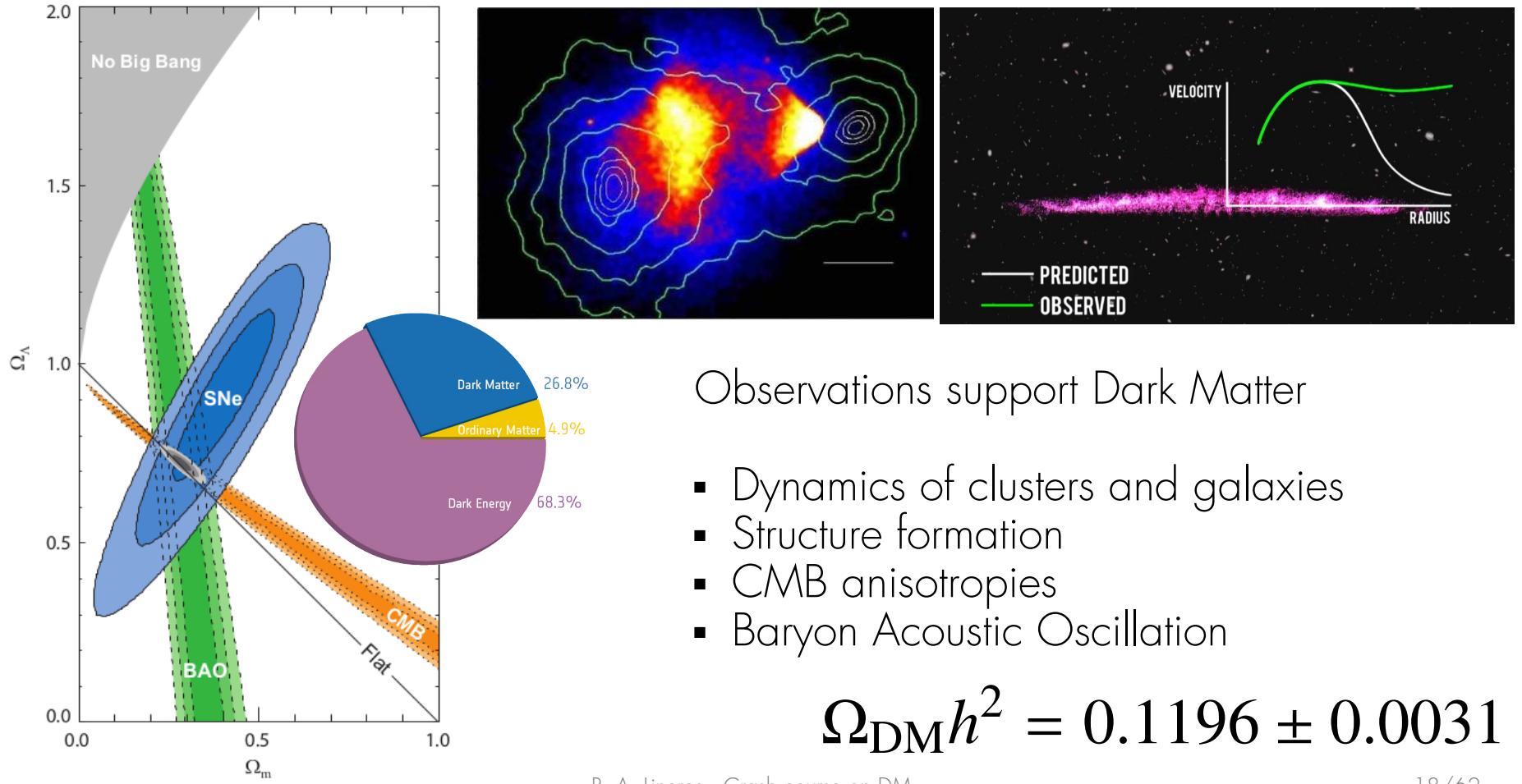
(short story) Dark Matter is everywhere!

Cosmic pie





Dark Matter



Galactic scales

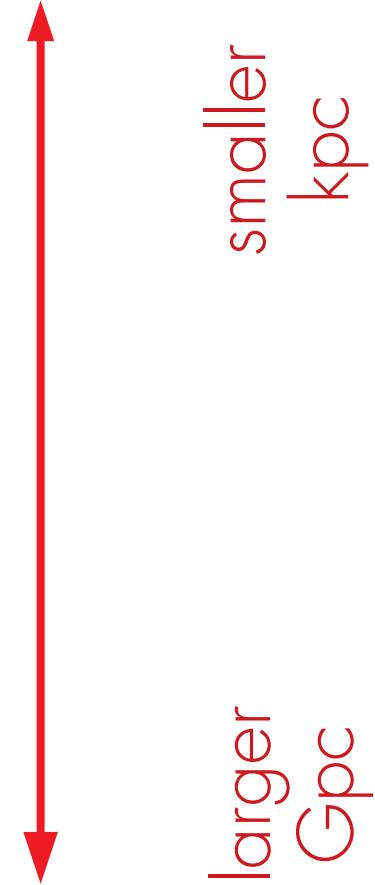
- Rotation curve
- Weak lensing
- Velocity dispersion of satellite galaxies
- Velocity dispersion of dSphs

Galaxy cluster scales

- Velocity dispersion of individual galaxies
- Strong and weak lensing
- Peculiar velocity flows
- X-ray emission

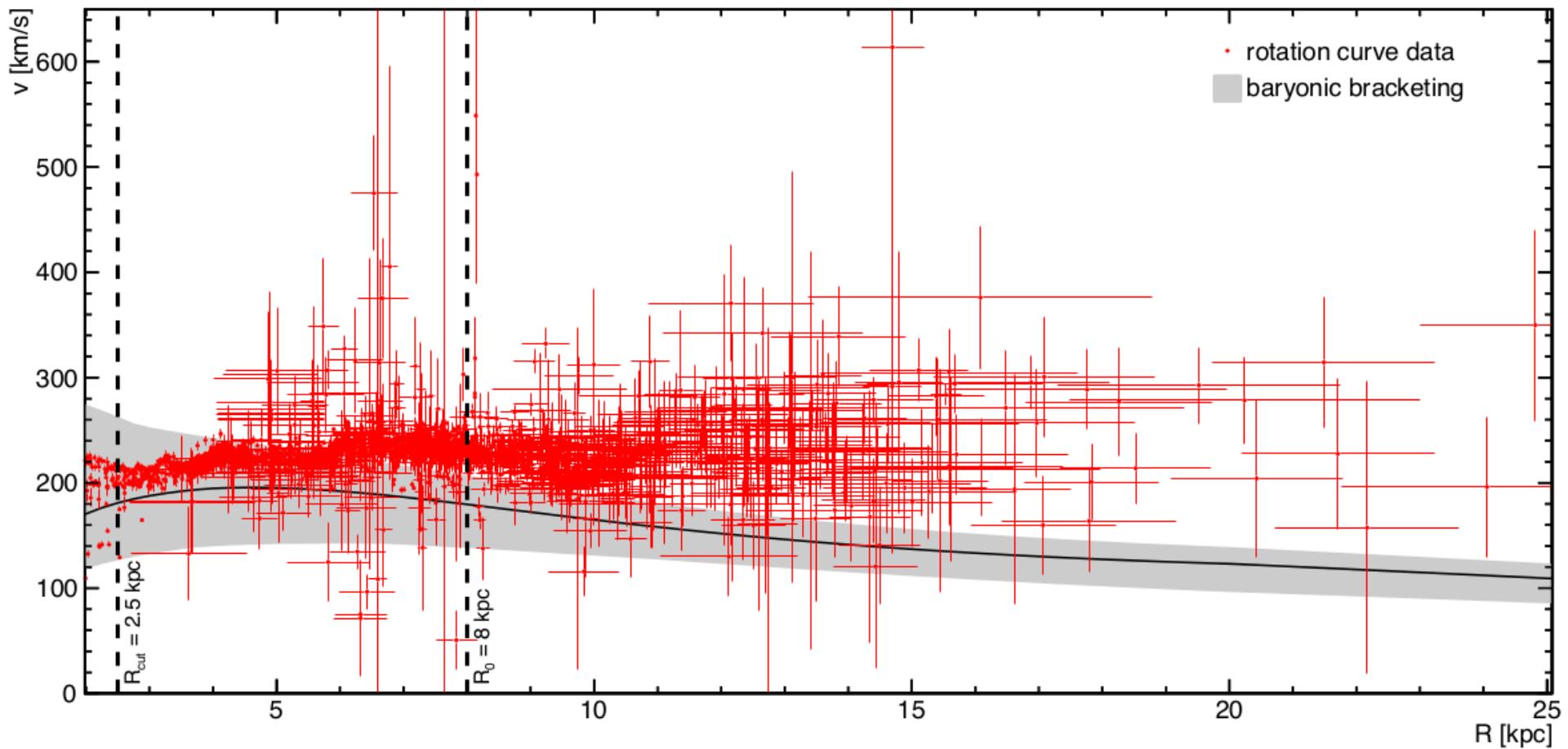
Cosmological scales

- CMB anisotropies
- Growth of structure
- LSS distribution
- BAOs
- SZ effect



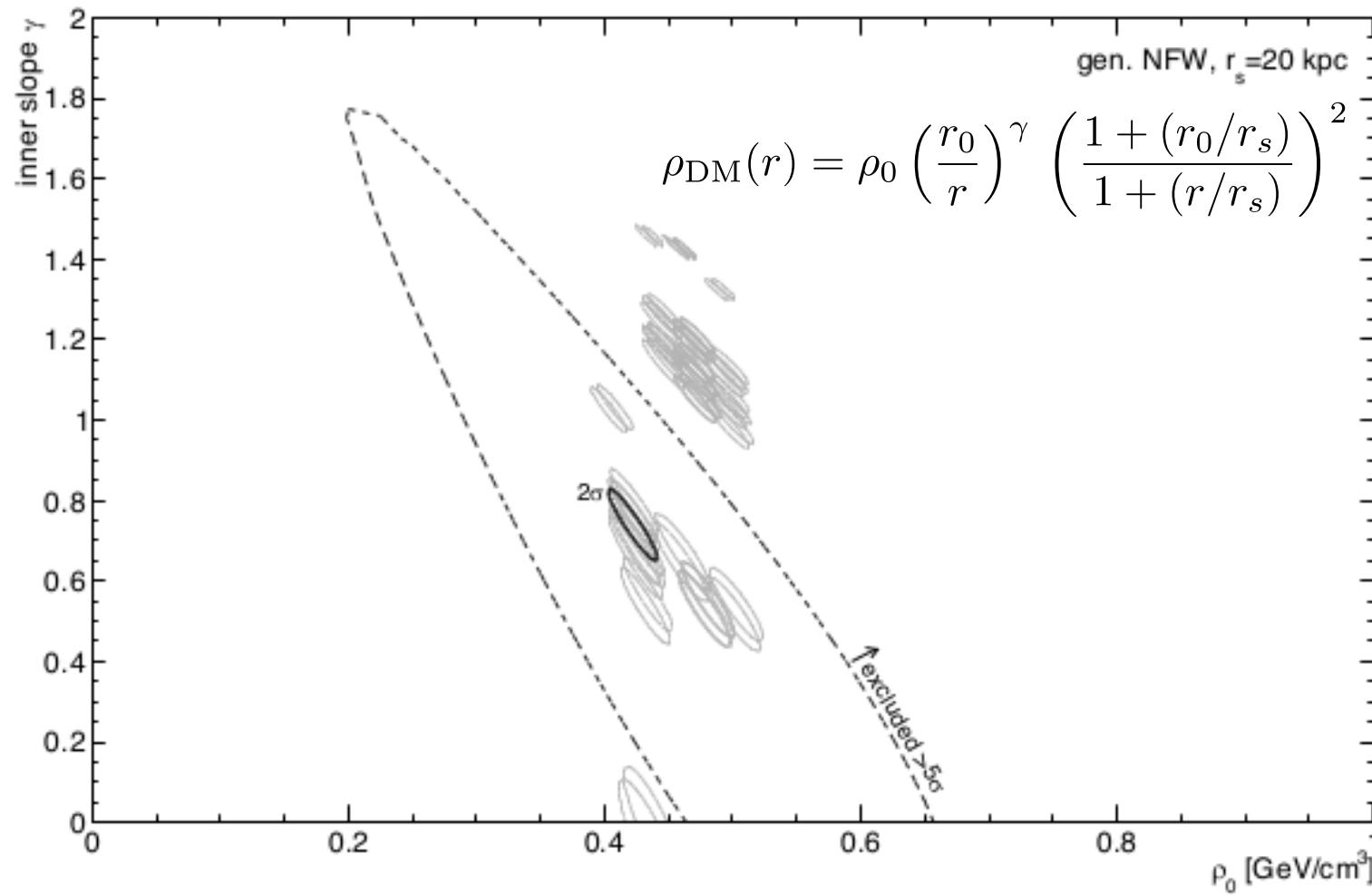
Galactic scales

M. Pato et al JCAP 1512 (2015)



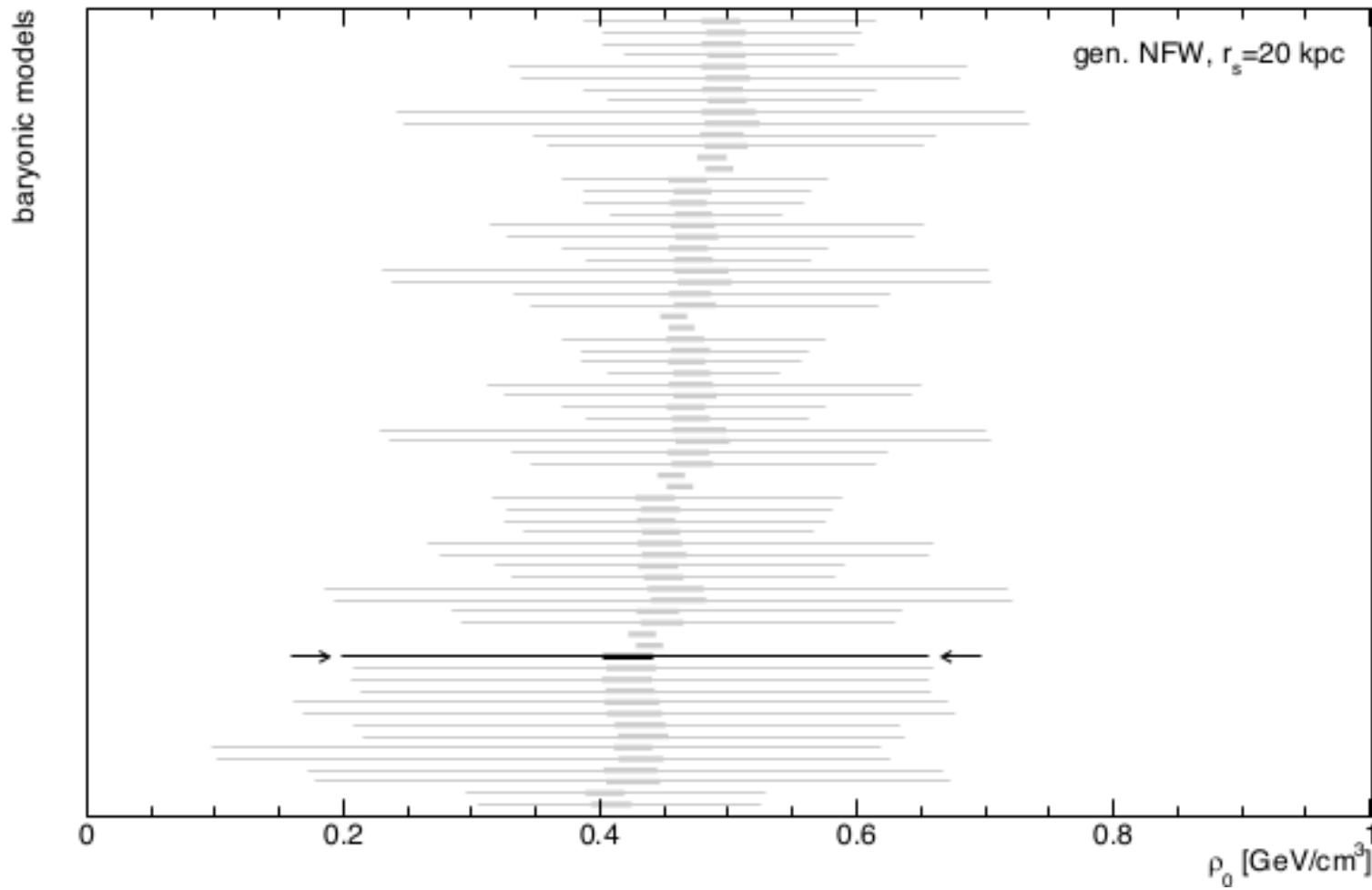
Galactic scale

M. Pato et al JCAP 1512 (2015)



Galactic scale

M. Pato et al JCAP 1512 (2015)



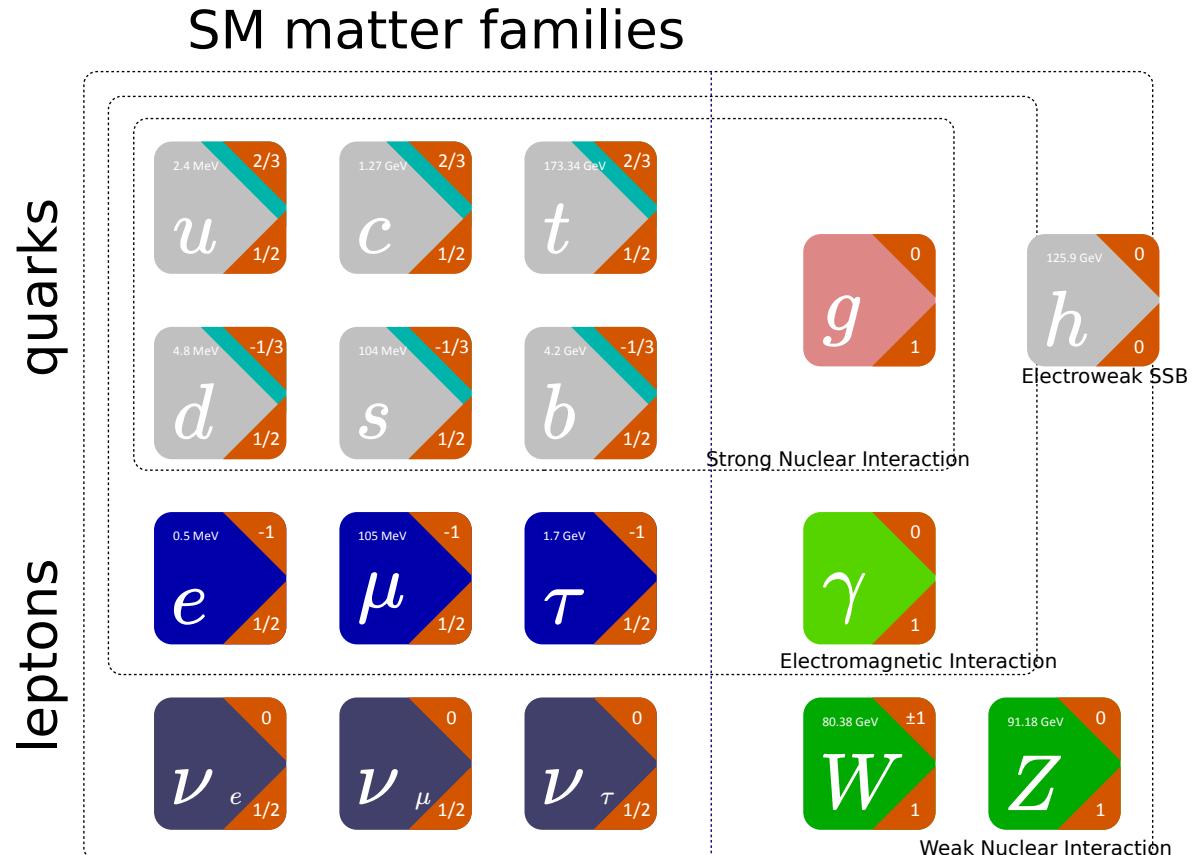


Dark Matter Candidates

(a small sample)

The Standard Model

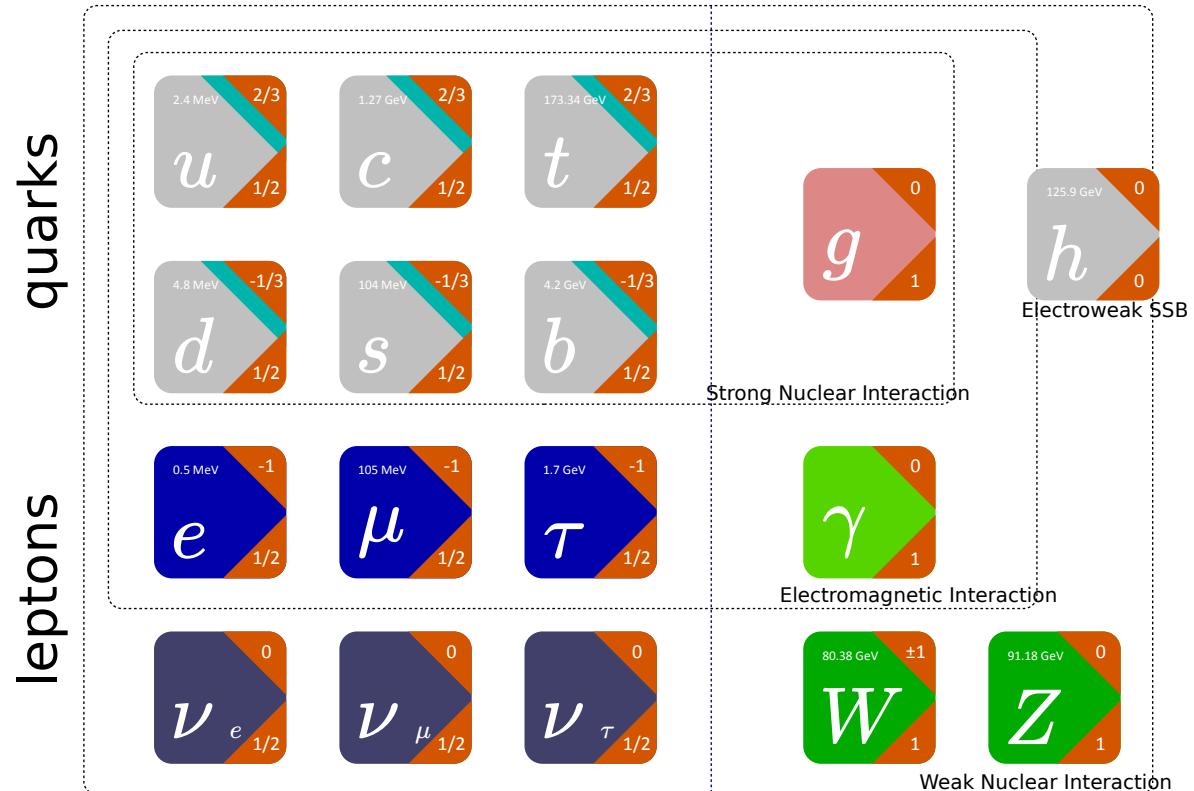
(so far)



The Standard Model

(so far)

SM matter families



Beyond SM



Dark Matter particle properties

Massive

Non baryonic

Electrically neutral

Stable

Dark Matter particle properties

Massive

Non baryonic

Electrically neutral

Stable



Dark Matter particle properties

Massive (*)

Non baryonic

Electrically neutral (**)

Stable (***)

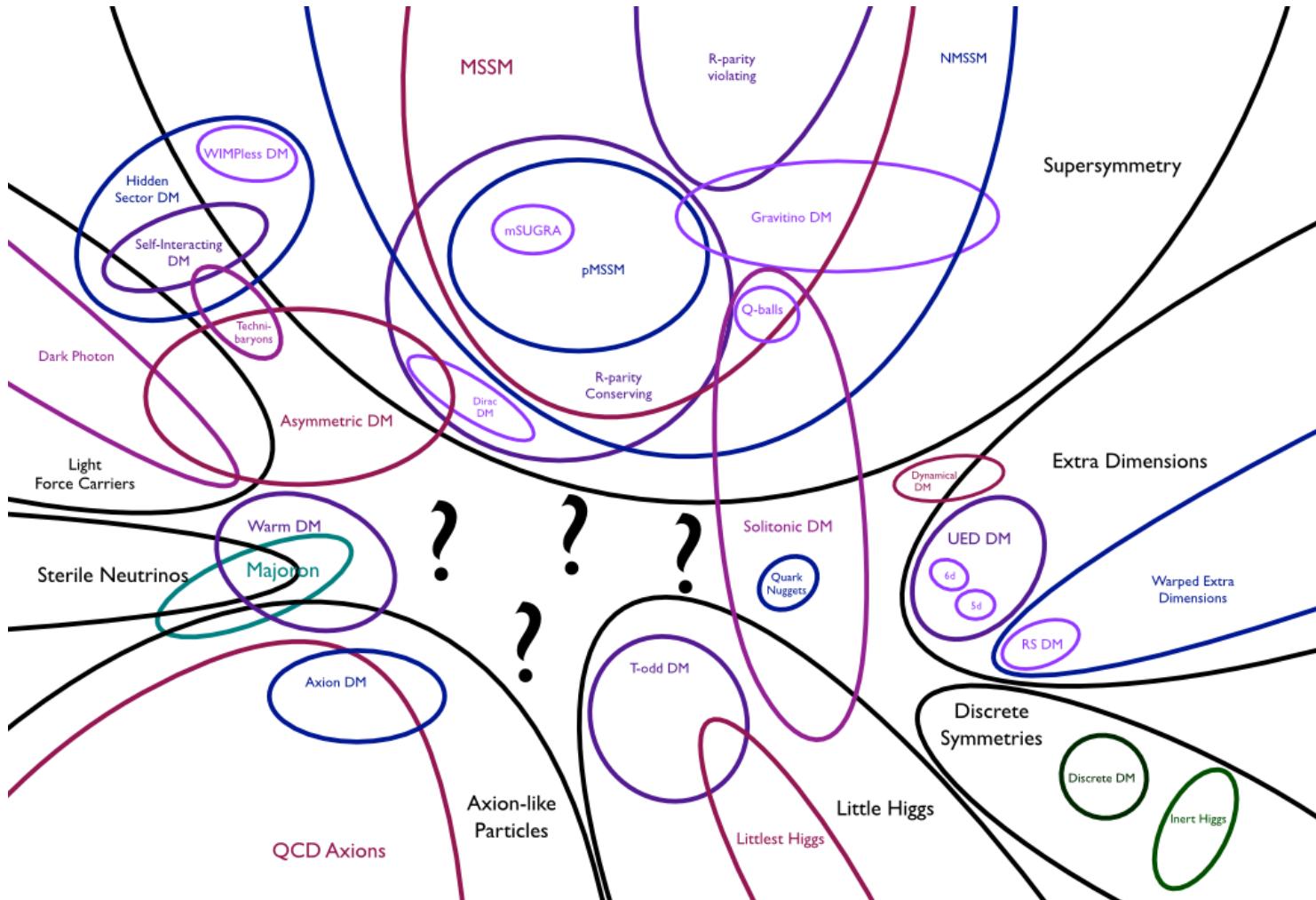
(*) Its mass can go from 10^{-22} eV to 10^9 GeV

(**) Except Milicharged DM or CHAMPs

(***) DM lifetime larger than 10^{27} seconds (Universe = 10^{17} seconds)

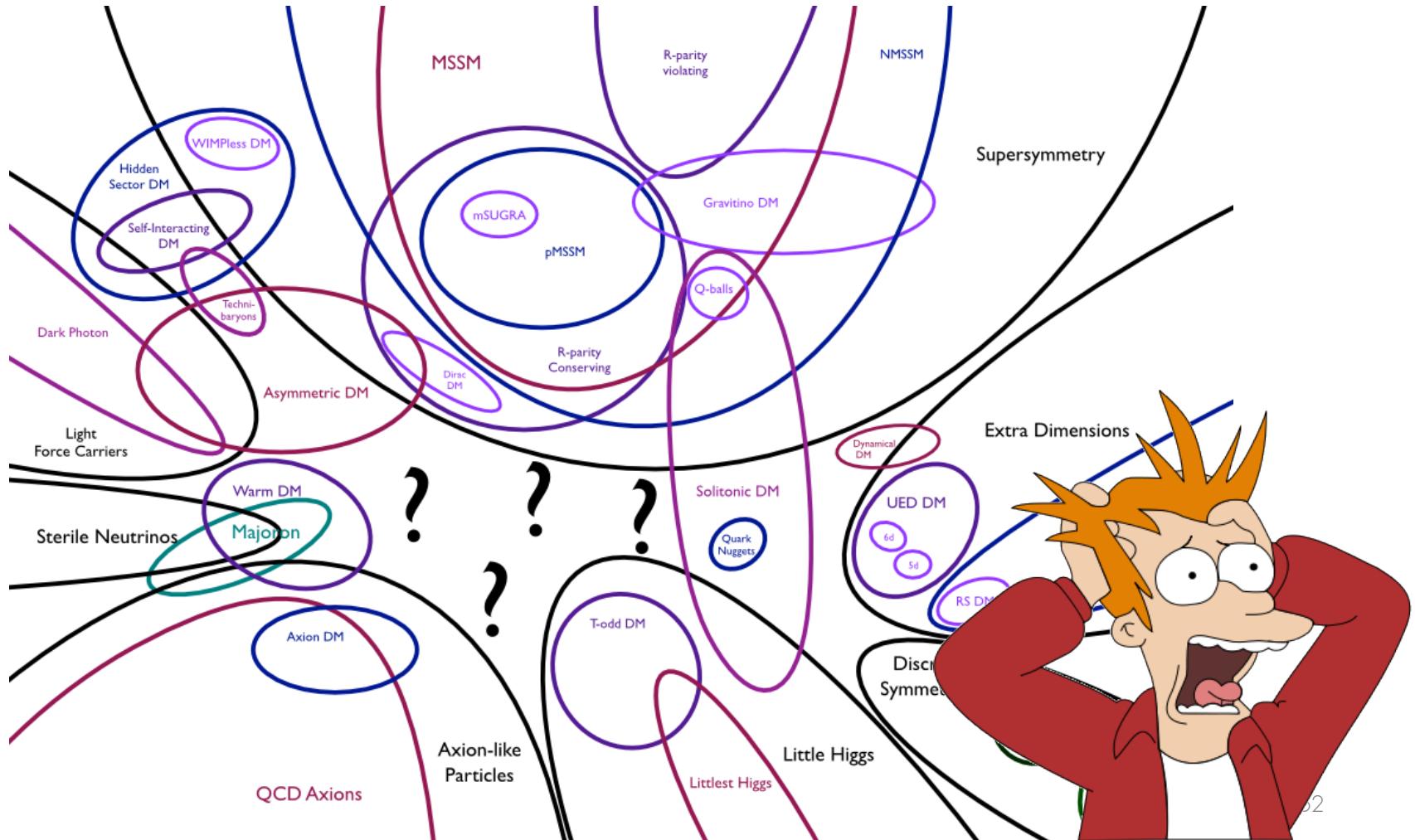


Some candidates



Adapted from 1401.6085

Some candidates



Adapted from 1401.6085



VS

Sterile Neutrinos

Sterile neutrinos

If right-handed neutrinos are part of the SM.
They are natural SM gauge singlets

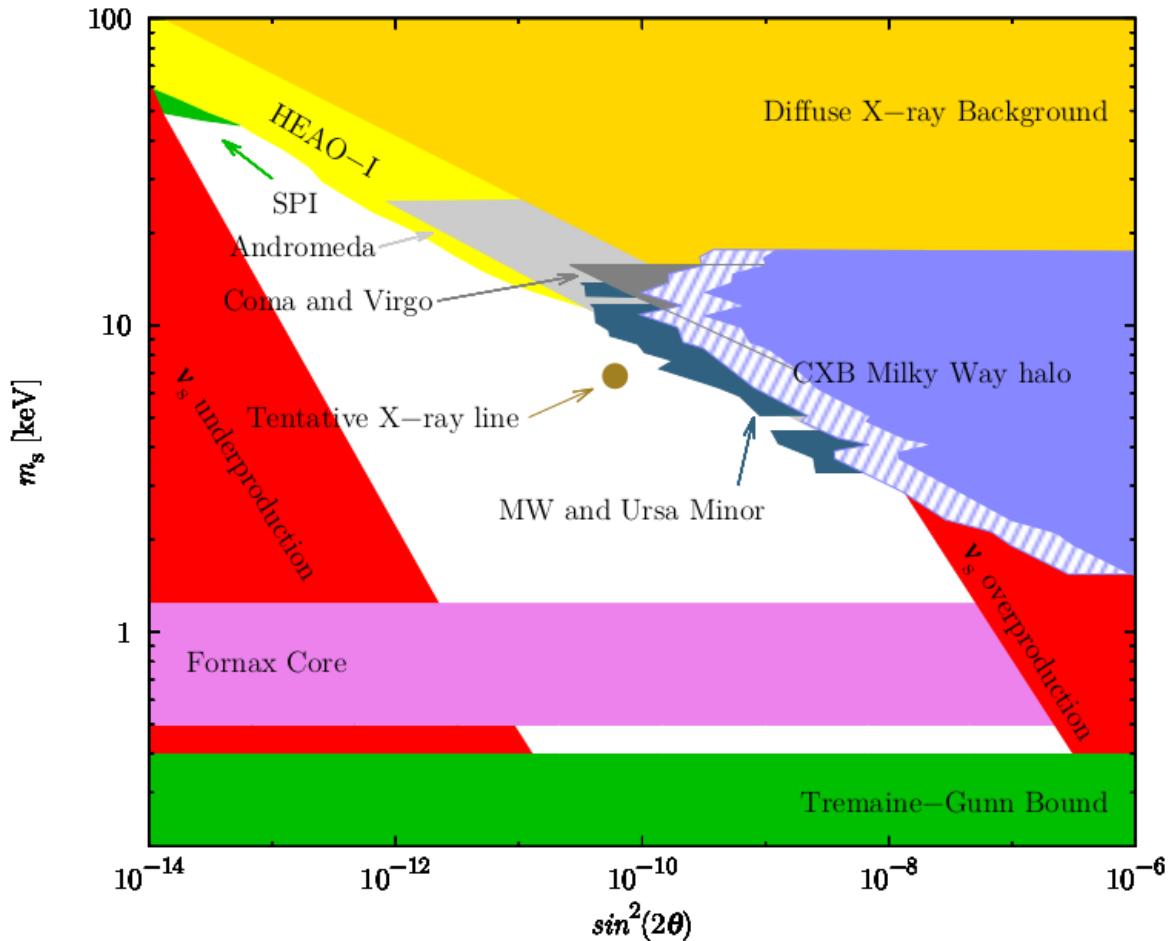
$$\mathcal{L}_2 = Y_2 \bar{L} H \nu_R$$

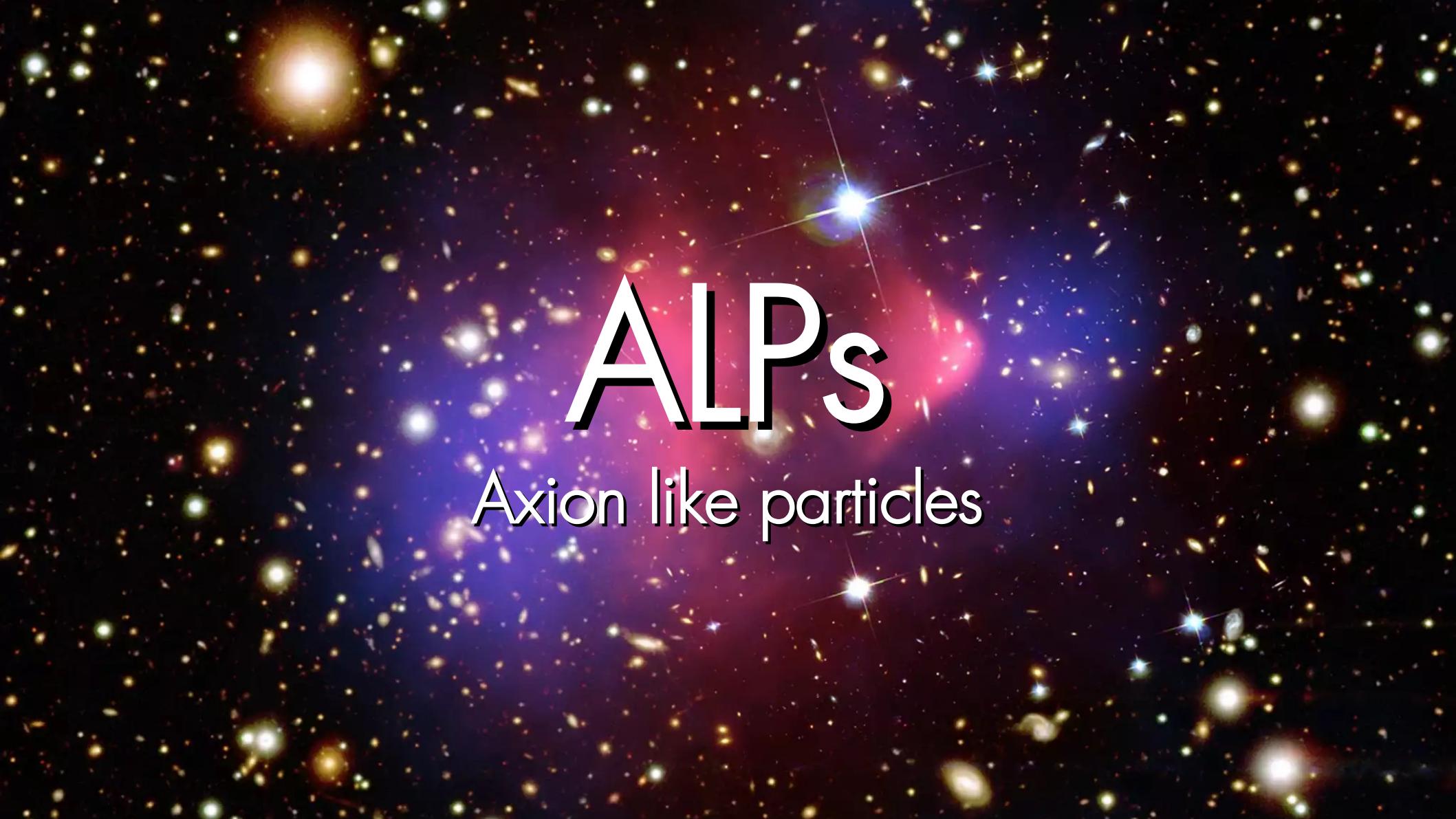
$$\mathcal{L}_R = \frac{1}{2} M_R \nu_R^T \nu_R + \text{h.c.}$$

$$\mathcal{M}_\nu = \begin{pmatrix} 0 & m_D \\ m_D & M_R \end{pmatrix}$$

Neutrino masses can be generated via the seesaw mechanism

Sterile neutrinos



The background of the image is a deep space scene filled with numerous galaxies and stars of varying sizes and colors, primarily in shades of blue, purple, and yellow.

ALPs

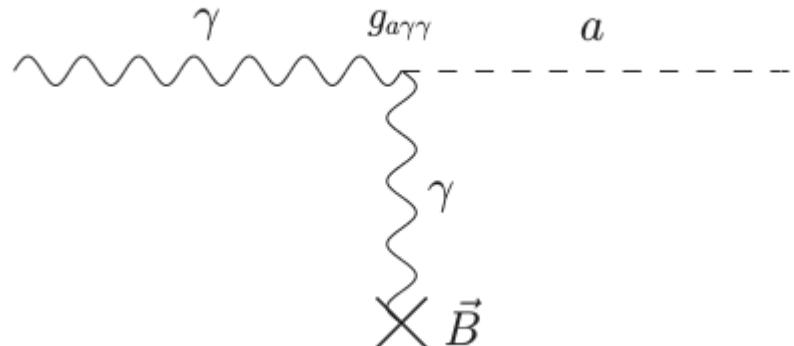
Axion like particles

Axion Like Particles

Pseudo-Goldstone bosons.

$$\mathcal{L} = -\frac{1}{4}g_\phi\phi F_{\mu\nu}\tilde{F}^{\mu\nu}$$

Its main feature is the coupling to photons



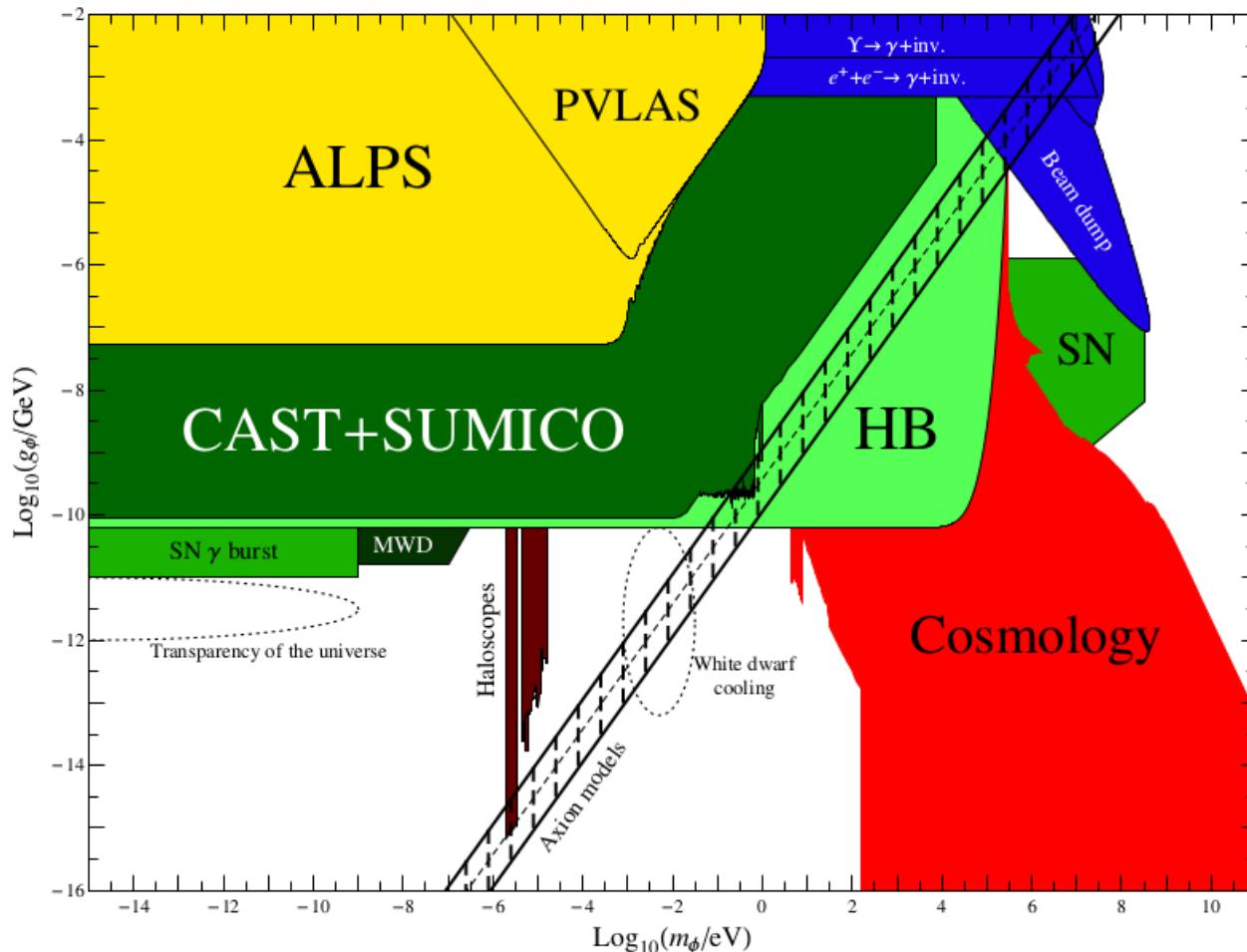
Primakoff effect

Searches involves photon emission line

Attenuation due to photon-axion conversion

See: 1210.3196

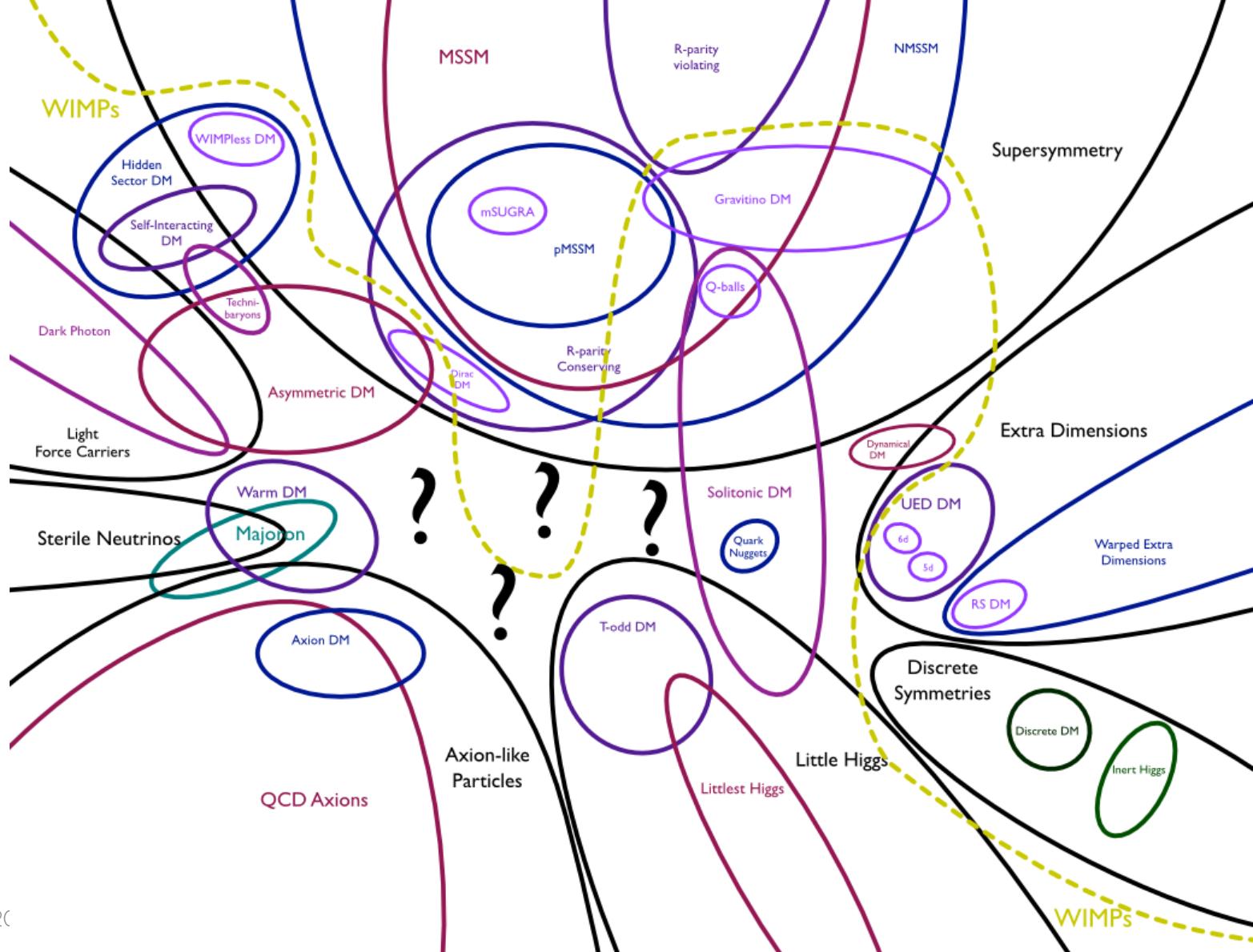
Axion Like Particles



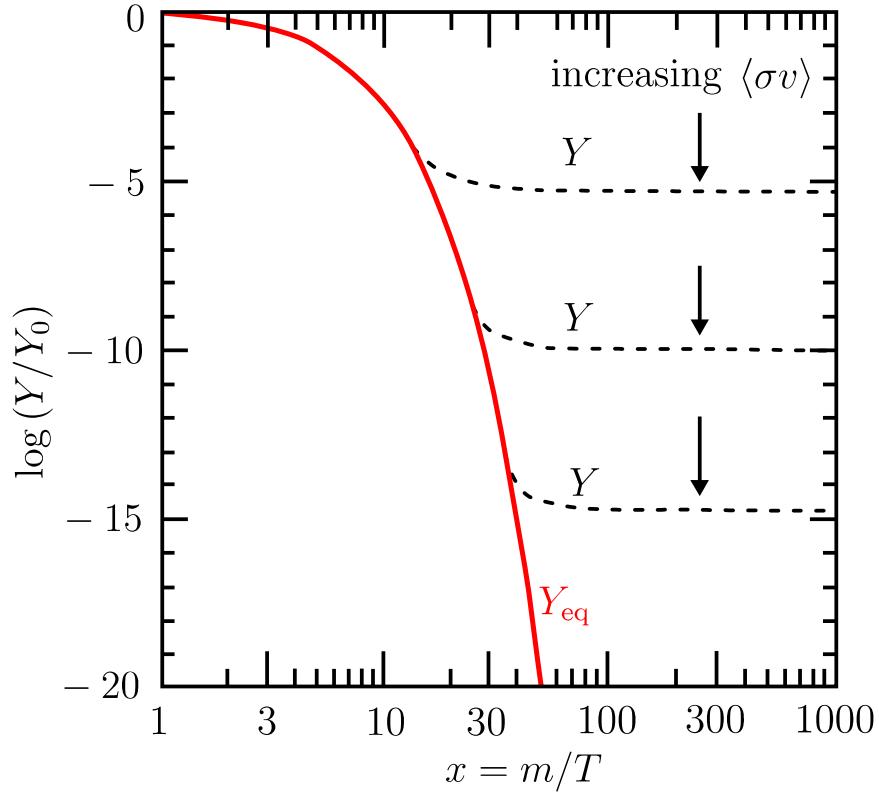


WIMPs

Weakly Interactive Massive Particles



WIMPs



Big Bang **Thermal relic**

Correct relic abundance for
 $\langle \sigma v \rangle \sim 1 \text{ pb} \cdot \text{c}$

Mass in **GeV-TeV** range

For WIMPs:

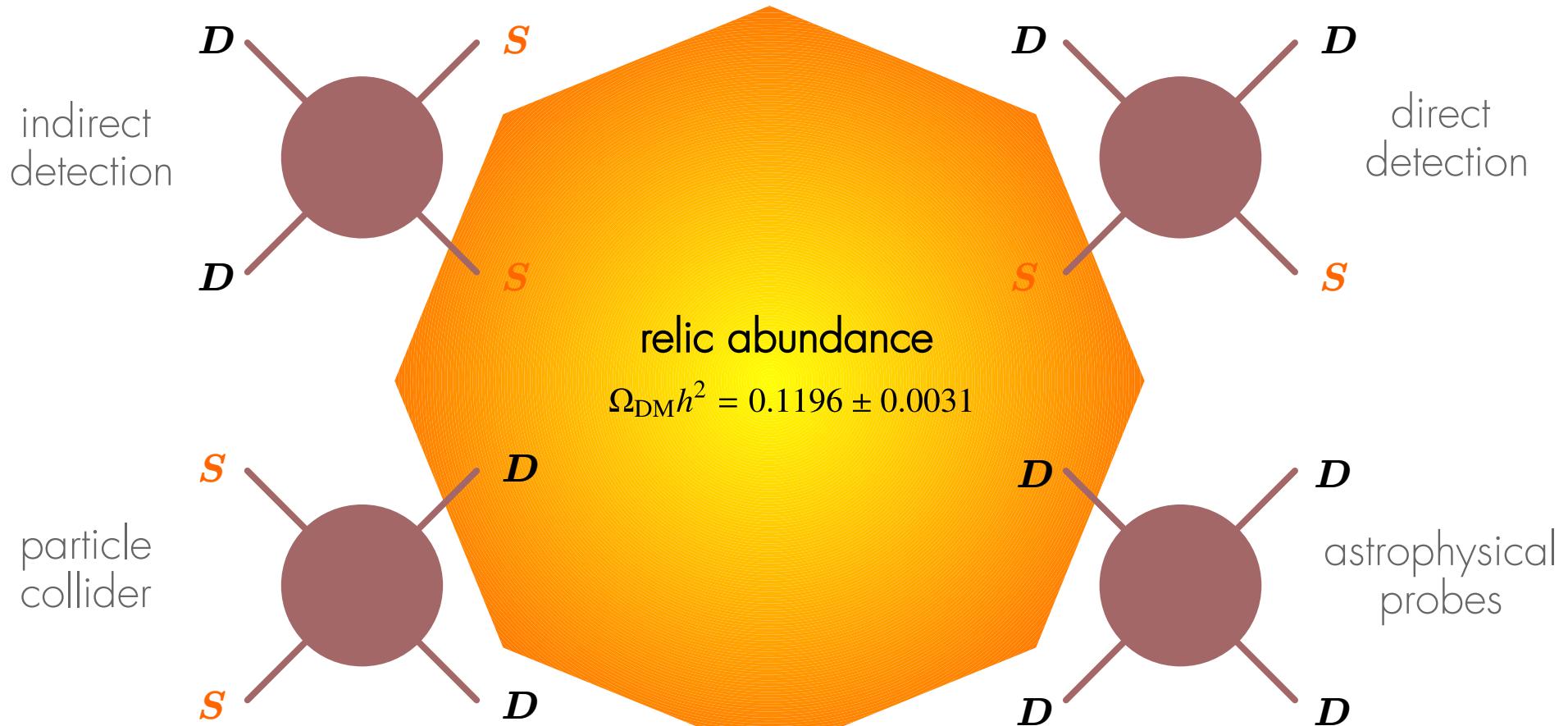
$$\Omega_{\text{DM}} h^2 \simeq 0.1 \frac{3 \times 10^{-26} \text{ cm}^3/\text{s}}{\langle \sigma v \rangle_{\text{f.o.}}}$$

8 Nov 2017

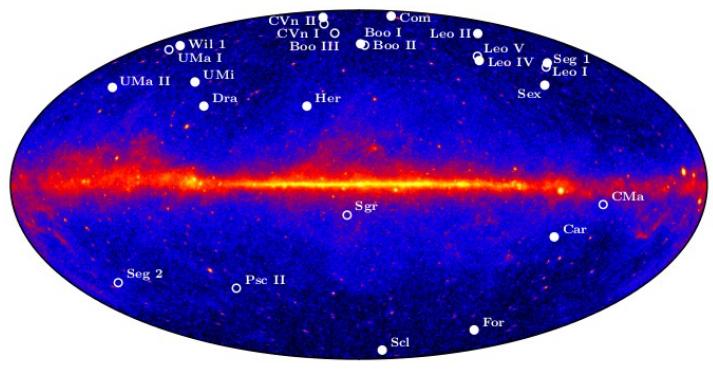
R. A. Lineros. Crash course on DM

$$T_{\text{DM}}^{\text{f.o.}} \simeq \frac{1}{20} m_{\text{DM}}$$

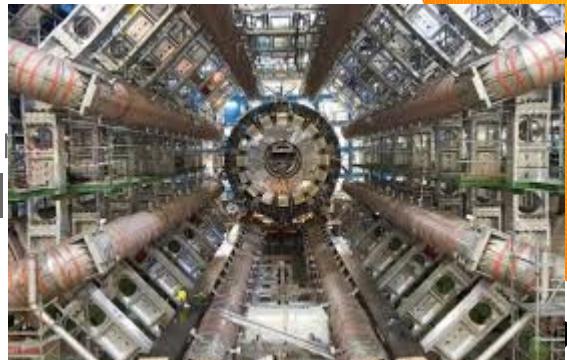
Dark Matter Searches



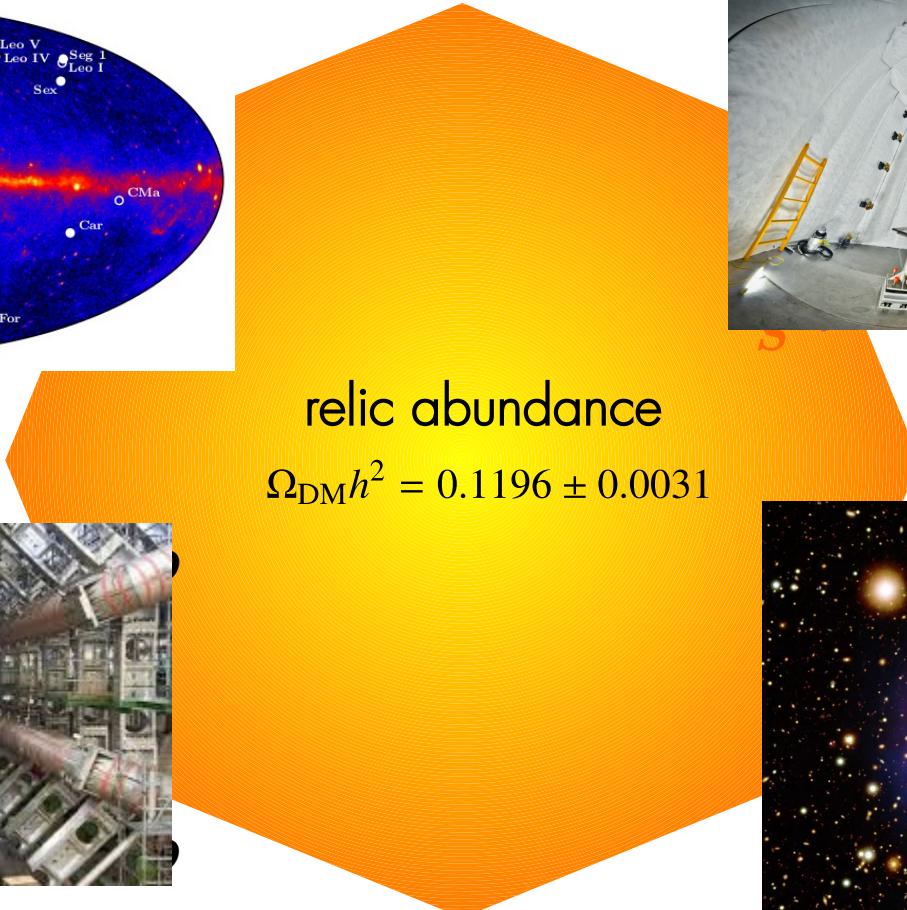
Dark Matter Searches



pa
co



8 Nov 2017



R A Lineros Crash course on DM

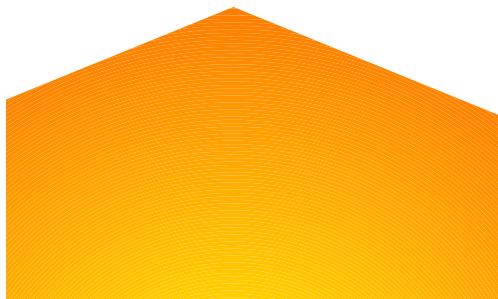
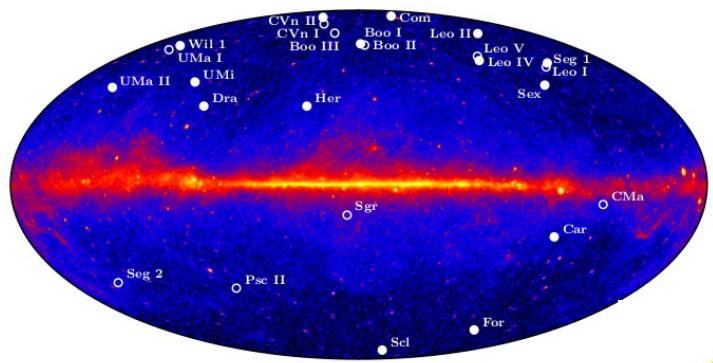


ct
tion

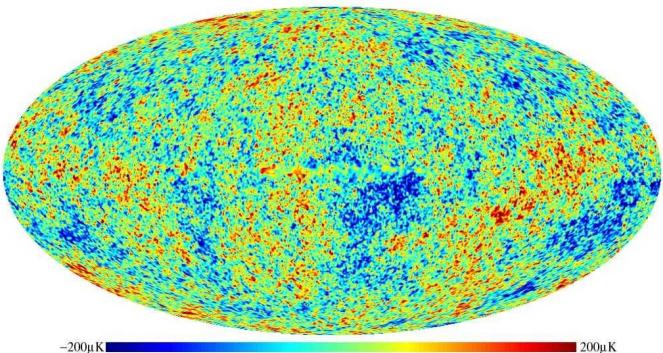
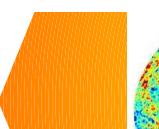


ical
S

Dark Matter Searches



ct
ction



par
col



Some extra details:



<https://youtu.be/DHc8Z2b1W5M>



<https://youtu.be/Gpi4vlQM348>



<https://youtu.be/Mxt33mN7sgU>



https://youtu.be/y-dpl_FulQY

Particle models for DM

$F_d = -bv$ $x(t) = X_m e^{-\frac{bt}{m}} \cos(\omega_n t + \phi)$ $\omega = \sqrt{\frac{k}{m} - \frac{b^2}{4m^2}}$ $E(t) \approx \frac{1}{2} k X_m^2 e^{-\frac{bt}{m}}$ $k(t) = \frac{1}{2} m v^2 = \frac{1}{2} m \omega^2$
 $T = 2\pi \sqrt{\frac{m}{k}}$ simple pendulum $\omega_d = \omega$ resonance
 $T = 2\pi \sqrt{\frac{m}{k+b^2/m}}$ physical period
 $v(t) = -\omega X_m \sin(\omega_n t + \phi)$
 $x(t) = X_m \cos(\omega_n t + \phi)$ $a(t) = -\omega^2 x(t)$
 $\gamma(x(t)) = Y_m \sin(kx - \omega t)$ wave in pos direction
 $k = \frac{2\pi}{\lambda}$ $V = \frac{\omega}{k} = \frac{\lambda}{T} = \lambda f$ $V = \sqrt{\frac{T}{\mu}}$ $\mu = \frac{mass}{length}$ $P_{ave} = \frac{1}{2} V k \omega^2 Y_m^2$
 resonance $\lambda = \frac{2L}{n}$ $n=1,2,3,\dots$ $V = \sqrt{\frac{\rho}{\rho_0}} \cdot \text{bulk modulus}$
 $f = \frac{V}{\lambda} = \frac{\Delta V}{2L}$ $n=1,2,3,\dots$ $P_m = 2PV$
 $P_m = V \rho \omega S_m$
 $I = \frac{\text{Power}}{\text{Area}} = \frac{P_s}{A}$
 pipe 2 open ends displacement anti-node pressure node
 pipe 1 open end displacement anti-node

Interference

$$\frac{\Delta L}{\lambda} = 0.5, 1, 2 \text{ fully constructive}$$

$$\frac{\Delta L}{\lambda} = 0.5, 1, 2, 3, 4 \text{ fully destructive}$$

$$Q = \frac{\Delta L}{\lambda} 2\pi r$$

$$B = 3\alpha$$

$$T_F = \frac{9}{5} T_C + 32$$

$$\Delta E_{int} = Q_{in} - Q_{out}$$

$$P_{cond} = \frac{Q}{t} = kA \frac{T_h - T_c}{L}$$

$$I = \frac{1}{2} \rho V w^2 S_m^2$$

$$Q = C_m \Delta T$$

$$f' = f \frac{V_d V_s}{V_s^2 + V_d^2}$$

$$Q = L m$$

$$W = \int_{p_0}^{p_f} V_r dp$$

$$\sin \theta = \frac{V_s}{V}$$

$$\frac{V_s}{V} = \text{mach } \#$$

$$B = (10) \log \frac{I}{I_0}$$

$$\log \frac{x}{y} = \log x - \log y$$

$$I_0 = 1 \times 10^{-12}$$

Particle model

$n = \frac{\text{molecules}}{6.02 \times 10^{23}}$

$R = \frac{k}{N_A}$ Multi-Slab
 $P_{\text{cond}} = \frac{A(T_u - T_c)}{\sum \frac{L}{k}}$ $P_{\text{radiation}} = \sigma \epsilon A T^4$ $P_{\text{net}} = P_{\text{abs}} - P_{\text{rad}}$

Adiabatic $Q=0 \quad \Delta E = -W$ $PV = nRT$ $R = 8.31 \sum \frac{1}{molar \cdot k}$ $5.67 \times 10^{-8} \frac{W}{m^2 K^4}$ $P_{\text{abs}} = \sigma \epsilon A T_{\text{env}}^4$

in Vol $W=0 \quad \Delta E = Q$

closed cyc $\Delta E = 0 \quad Q = W$ $W = A R T \ln \frac{V_f}{V_i} (\text{isothermal})$ $k_{\text{Boltz}} = \frac{3}{2} kT$ $M = \text{molar mass}$ $V_{\text{rms}} = \sqrt{\frac{3RT}{M}}$

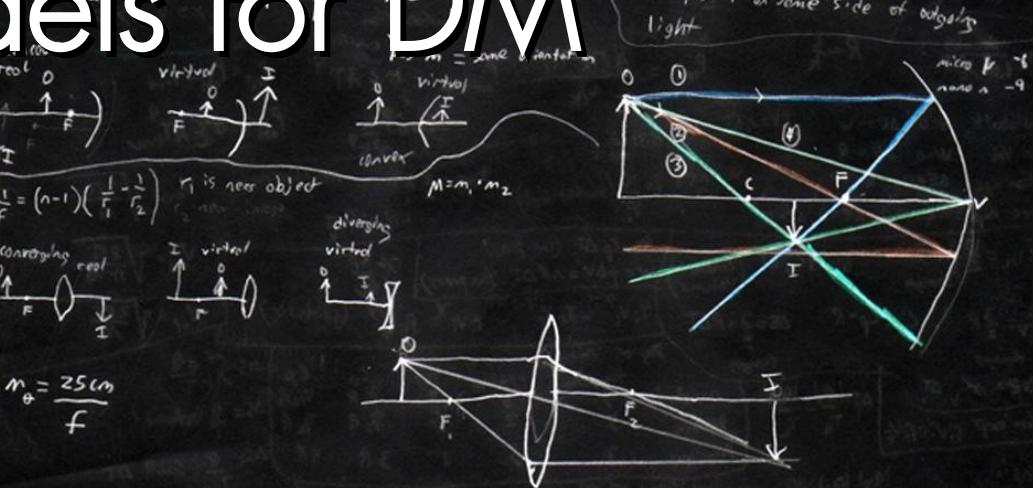
Free exp $Q = W = 0 \quad \Delta E = 0 \quad \Delta T = 0$ $Q = \sigma \epsilon A T (\text{iso. area})$ $1.38 \times 10^{-23} \frac{J}{K}$

$$\begin{aligned}
 E_{int} &= \frac{3}{2} \gamma R T & Q = \gamma C_V \Delta T & \boxed{\text{constant Volume}} & R = \gamma C_P \Delta T \quad (\text{constant pressure}) \\
 \text{changes} & & W = P \Delta V = \gamma R \Delta T & & \boxed{\text{constant pressure}} & V_{avg} = \sqrt{\frac{8 \pi T}{\gamma R M}} \\
 C_V &= \frac{3}{2} R & C_V &= C_P - R & \gamma = \frac{C_P}{C_V} & V_{mp} = \sqrt{\frac{2 \pi T}{M}} \\
 \text{changes} & & \Delta E = \frac{3}{2} \gamma R \Delta T & & T_1 V_1^{\gamma-1} = T_2 V_2^{\gamma-1} & \\
 \Delta E &= \gamma C_V \Delta T & \text{degrees of freedom} & & P_1 V_1^{\gamma} = P_2 V_2^{\gamma} & \text{(adiabatic)} \\
 \text{Free exp } P_1 V_1 = P_2 V_2 & & \text{translational} & & \Delta S = \int \frac{dQ}{T} & \\
 T_H & & & & & \Delta S = \frac{Q}{T} \text{ isothermal} \\
 T_L & & & & & \\
 \text{Polytropic} & & & & & \Delta S = \frac{Q}{T} \text{ adiabatic} \\
 \text{Enthalpy} & & & & & \Delta S = \frac{Q}{T} \text{ DT}
 \end{aligned}$$

$$\text{we } |Q_{\text{H}}| - |Q_L| \quad \frac{|Q_H|}{T_H} = \frac{|Q_L|}{T_L} \quad E = \frac{|W|}{|Q_H|} \text{ energy we get} \quad E_C = \frac{|Q_{\text{H}}| - |Q_L|}{|Q_{\text{H}}|} = 1 - \frac{|Q_L|}{|Q_{\text{H}}|} \quad \Delta S = \eta R \ln \left(\frac{V_f}{V_i} \right) + n C_v \ln \frac{T_f}{T_i}$$

refrigerator $k = |Q_L|$ what we want
 $|W|$ what we pay for $E_C = \frac{|Q_L|}{|Q_{\text{H}}|}$

critical angle $n_1 \sin \theta_c = n_2 \sin \theta_2$
 $\theta_c = \sin^{-1} \frac{n_2}{n_1}$
 Brewster's angle $\theta_B = \tan^{-1} \frac{n_2}{n_1}$
 single slit diffraction $a \sin \theta = m\lambda$ ($m=1, 2, 3$) minima
 $I = I_m \left(\frac{\sin \theta}{\theta} \right)^2$ $\theta = \frac{1}{2} \alpha = \frac{\pi r}{\lambda} \sin \theta$
 diffraction grating $d \sin \theta = m\lambda$ ($m=0, 1, 2$) maxima 1 line
 $\frac{d}{\lambda} = \frac{m}{\sin \theta}$ dispersion
 radius r curvature $\frac{1}{r} + \frac{1}{f} = \frac{1}{f}$ spherical mirror
 $R = Nn'$ resolving power
 $\Delta \theta_{hw} = \frac{\lambda}{Nd \cos \theta}$ half width
 $\Delta L = d \sin \theta$ double slit
 $2L = \frac{m\lambda}{\sin \theta}$ ($m=0, 1, 2$) minima
 $\Delta x = \frac{L\lambda}{d}$ $\frac{1}{d} = \frac{\lambda}{\Delta x}$ with separation
 circular diffraction $\sin \theta = 1.22 \frac{\lambda}{d}$ first minimum
 aperture diameter
 Rayleigh criterion $\theta = 1.22 \frac{\lambda}{d}$ rayleigh criterion
 if $\theta > \theta_r$ it is resolved,
 Plane Mirrors $i = -p$
 Obj - S is pos if on same side as incoming light
 Img - S' is pos if on same side of outgoing light
 F - pos if on same side of all



Inert singlet DM model

The simplest DM model.

$$\mathcal{L} = \frac{1}{2} (\partial_\mu \phi)^2 - \frac{m_\phi}{2} \phi^2 - \frac{\lambda_\phi}{4} \phi^4 - \frac{\lambda_{\phi h}}{2} \phi^2 H^\dagger H + \mathcal{L}_{SM}$$

Features:

DM is a real scalar charged with a Z_2 symmetry

The interaction with the SM is via Higgs particle

The relevant parameter are mass and coupling to the Higgs

Inert singlet DM model

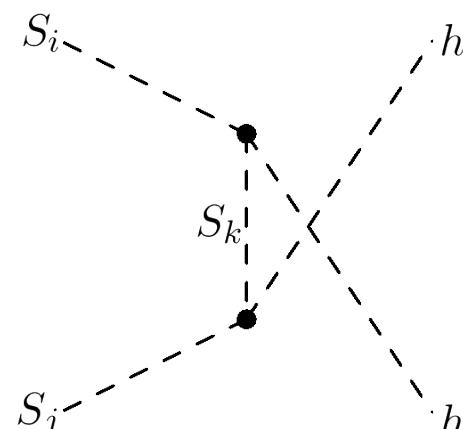
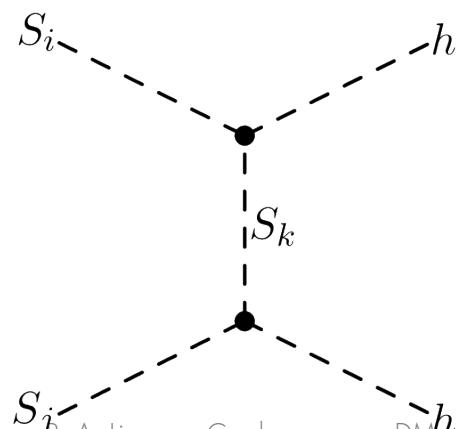
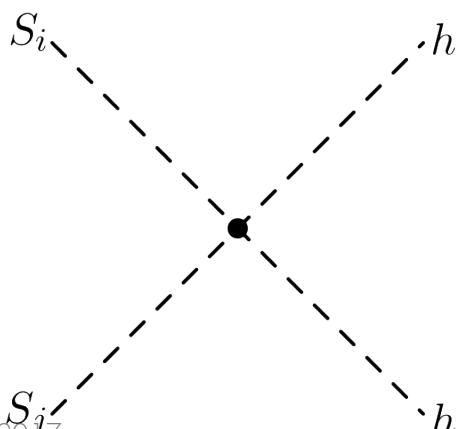
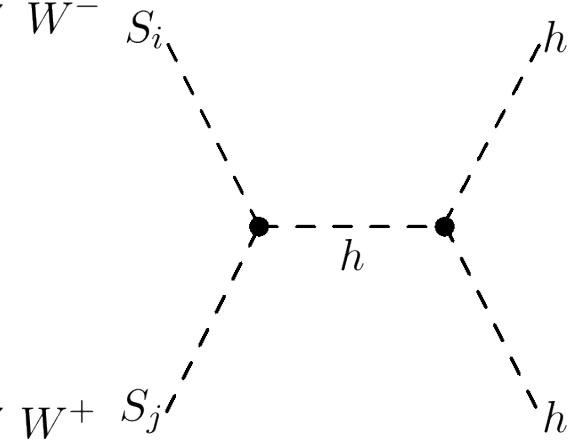
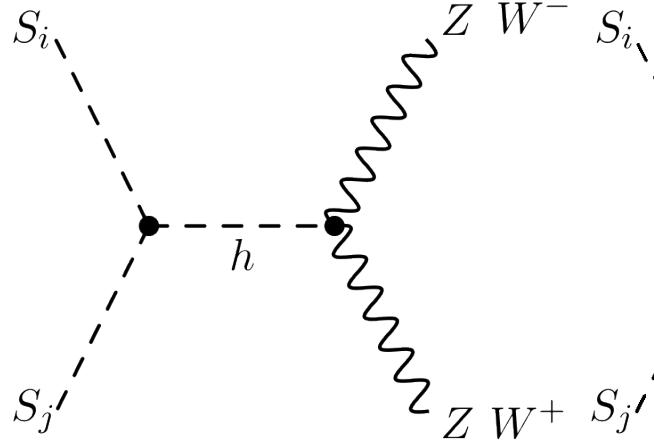
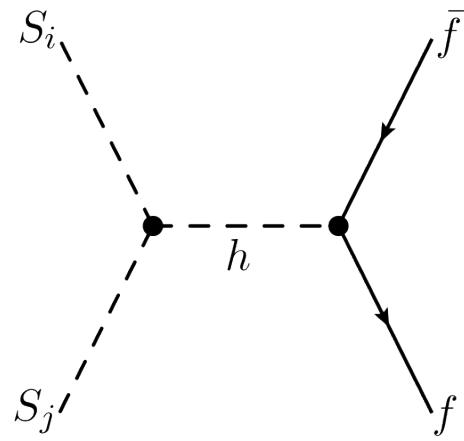
The simplest DM model.

$$\mathcal{L} = \frac{1}{2} (\partial_\mu \phi)^2 - \frac{m_\phi}{2} \phi^2 - \frac{\lambda_\phi}{4} \phi^4 - \frac{\lambda_{\phi h}}{2} \phi^2 H^\dagger H + \mathcal{L}_{SM}$$

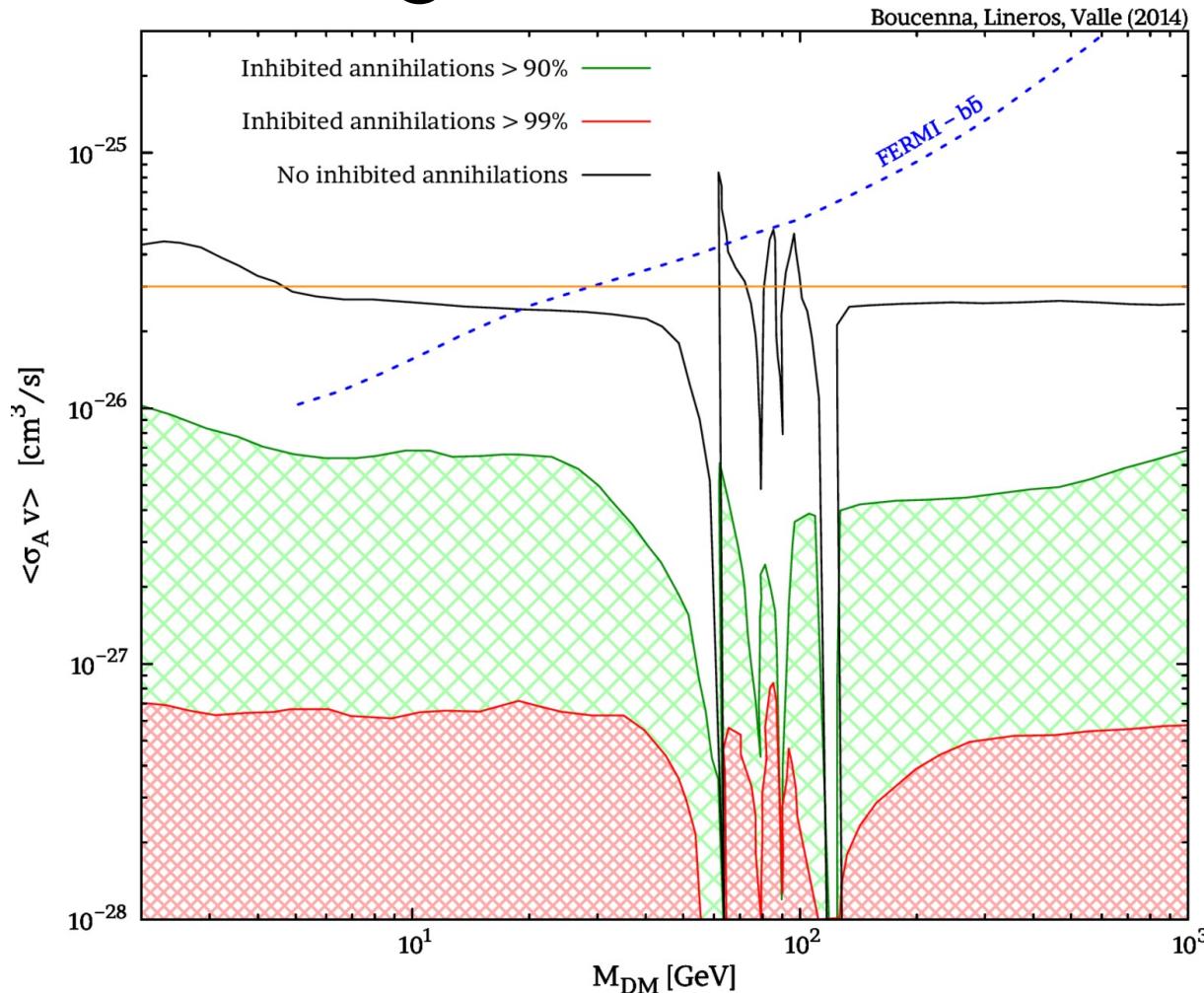
Small exercise:

Draw diagrams relevant for the relic abundance

Inert singlet DM model



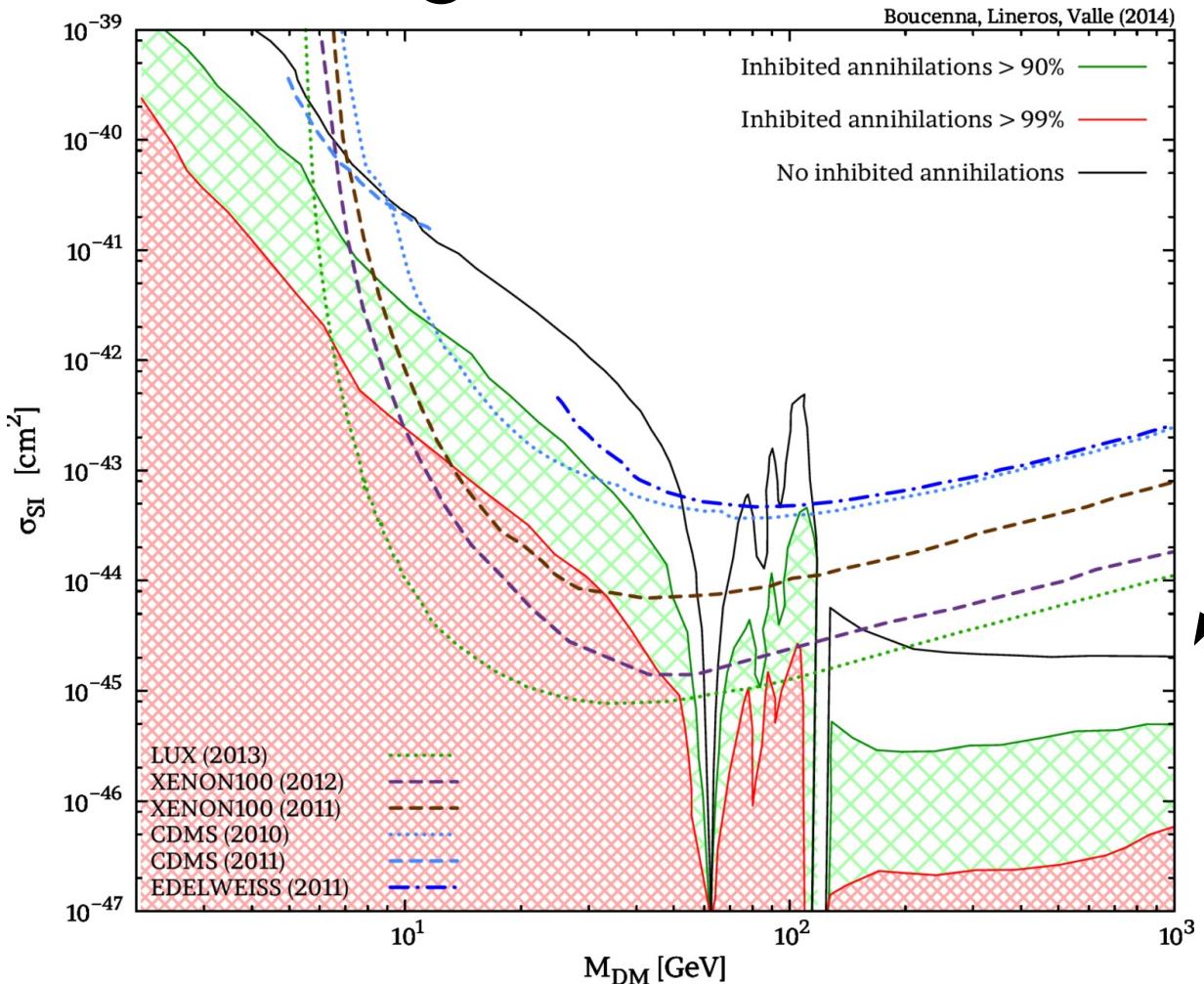
Inert singlet DM model



Solution is the black line

Tight relation among
mass and higgs coupling

Inert singlet DM model



Solution is the black line
Tight relation among
mass and higgs coupling

Fermion singlet DM model

Dark Matter is a majorana fermion connected to the SM via a scalar singlet

$$\begin{aligned}\mathcal{L} = & \frac{1}{2} \bar{\psi} (i\gamma_\mu \partial_\mu - m_\psi) \psi + \omega \phi \bar{\psi} \psi + \frac{1}{2} (\partial_\mu \phi)^2 - \frac{m_\phi^2}{2} \phi^2 \\ & - \frac{\lambda_1}{3} \phi^3 - \frac{\lambda_2}{4} \phi^4 - \frac{\lambda_3}{2} \phi H^\dagger H - \frac{\lambda_4}{2} \phi^2 H^\dagger H + \mathcal{L}_{SM}\end{aligned}$$

The model contains:

2 extra particles and 7
parameters

More freedom to reproduce DM
observables

Small exercise:

Draw diagrams
relevant for the relic
abundance

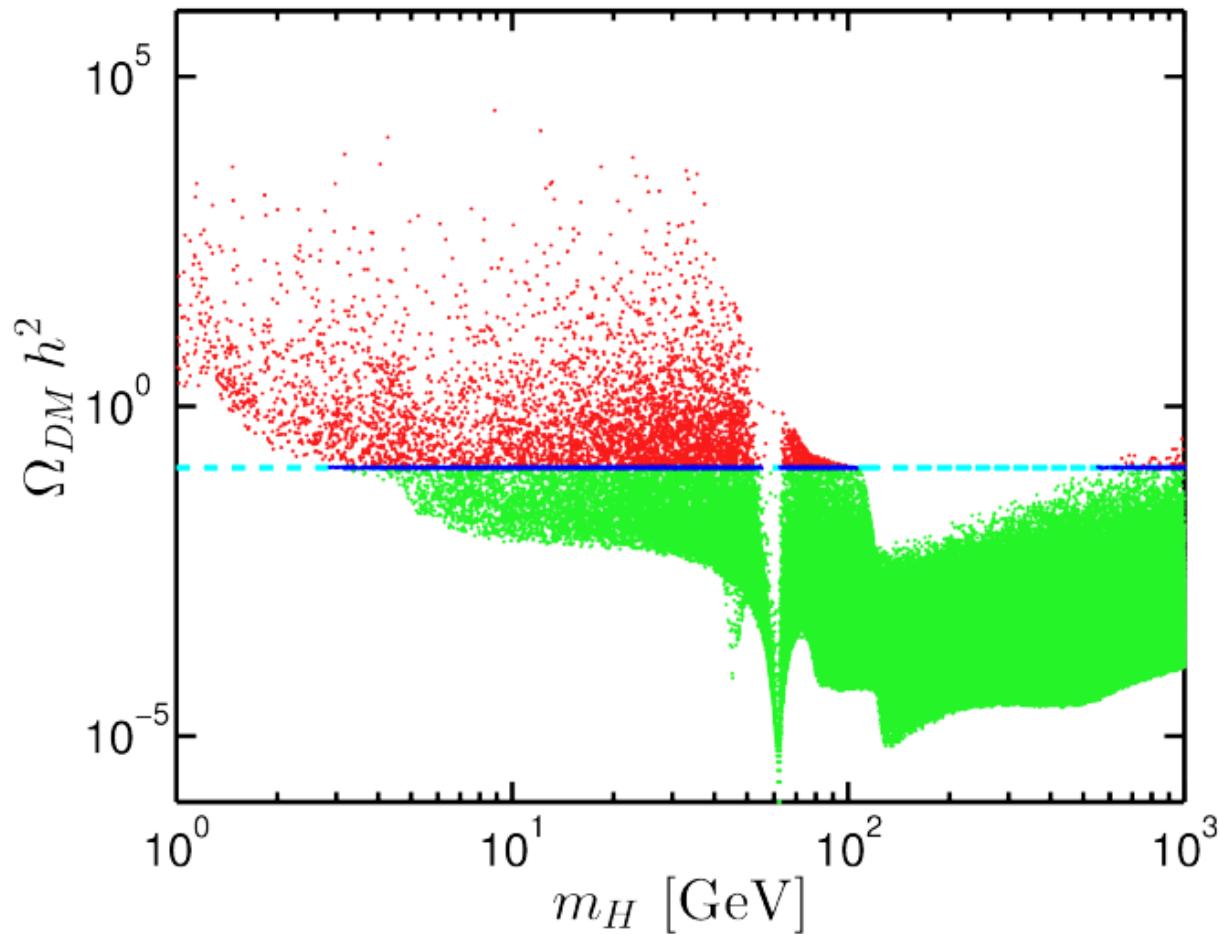
Inert Higgs DM model

DM is part of a SU(2) scalar doublet (copy of Higgs) but charged with a Z_2 symmetry

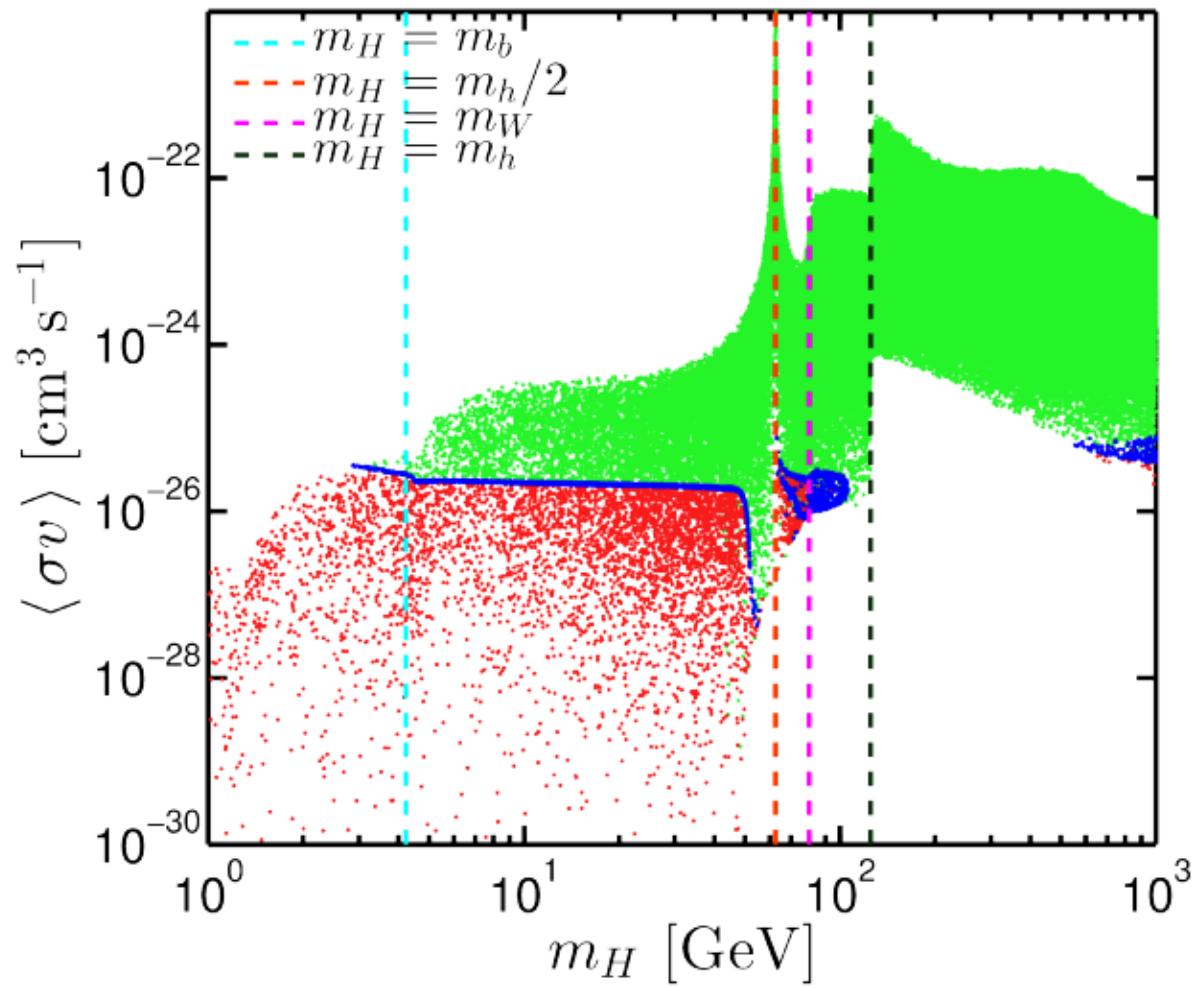
$$V = \mu_1^2 |H_1|^2 + \mu_2^2 |H_2|^2 + \lambda_1 |H_1|^4 + \lambda_2 |H_2|^4 + \lambda_3 |H_1|^2 |H_2|^2 + \lambda_4 |H_1^\dagger H_2|^2 + \frac{\lambda_5}{2} [(H_1^\dagger H_2)^2 + \text{h.c.}] ,$$

$$H_1 = \begin{pmatrix} G^+ \\ \frac{1}{\sqrt{2}} (v + h + iG^0) \end{pmatrix} , \quad H_2 = \begin{pmatrix} H^+ \\ \frac{1}{\sqrt{2}} (S + iA) \end{pmatrix}$$

Inert Higgs DM model



Inert Higgs DM model

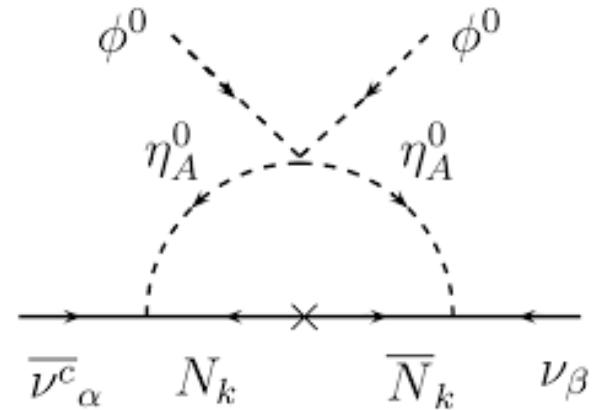


Scotogenic DM model

(Ma, 2006... and many other papers)

Model provides a mechanism where neutrino acquire mass via loop.

Inside the loop DM particles run.



$$(\mathcal{M}_\nu)_{ij} = \sum_k \frac{h_{ik} h_{jk} M_k}{16\pi^2} \left[\frac{m_R^2}{m_R^2 - M_k^2} \ln \frac{m_R^2}{M_k^2} - \frac{m_I^2}{m_I^2 - M_k^2} \ln \frac{m_I^2}{M_k^2} \right],$$

Scotogenic DM model

$$\mathcal{L} = \mathcal{L}_{\text{SM}} - \frac{1}{2} M_i \overline{N}_i N_i^{\mathcal{C}} + h_{ij} \overline{N}_i \tilde{\eta}^\dagger \ell^j + \text{h.c.}$$

$$\begin{aligned} V = & m_H^2 \phi^\dagger \phi + m_\eta^2 \eta^\dagger \eta + \frac{\lambda_1}{2} (\phi^\dagger \phi)^2 \\ & + \frac{\lambda_2}{2} (\eta^\dagger \eta)^2 + \lambda_3 (\phi^\dagger \phi) (\eta^\dagger \eta) \\ & + \lambda_4 (\phi^\dagger \eta) (\eta^\dagger \phi) + \left[\frac{\lambda_5}{2} (\eta^\dagger \phi)^2 + \text{h.c.} \right] \end{aligned}$$

Singlet-Triplet Scotogenic DM

DM is mix between a singlet fermion and a SU(2) triplet

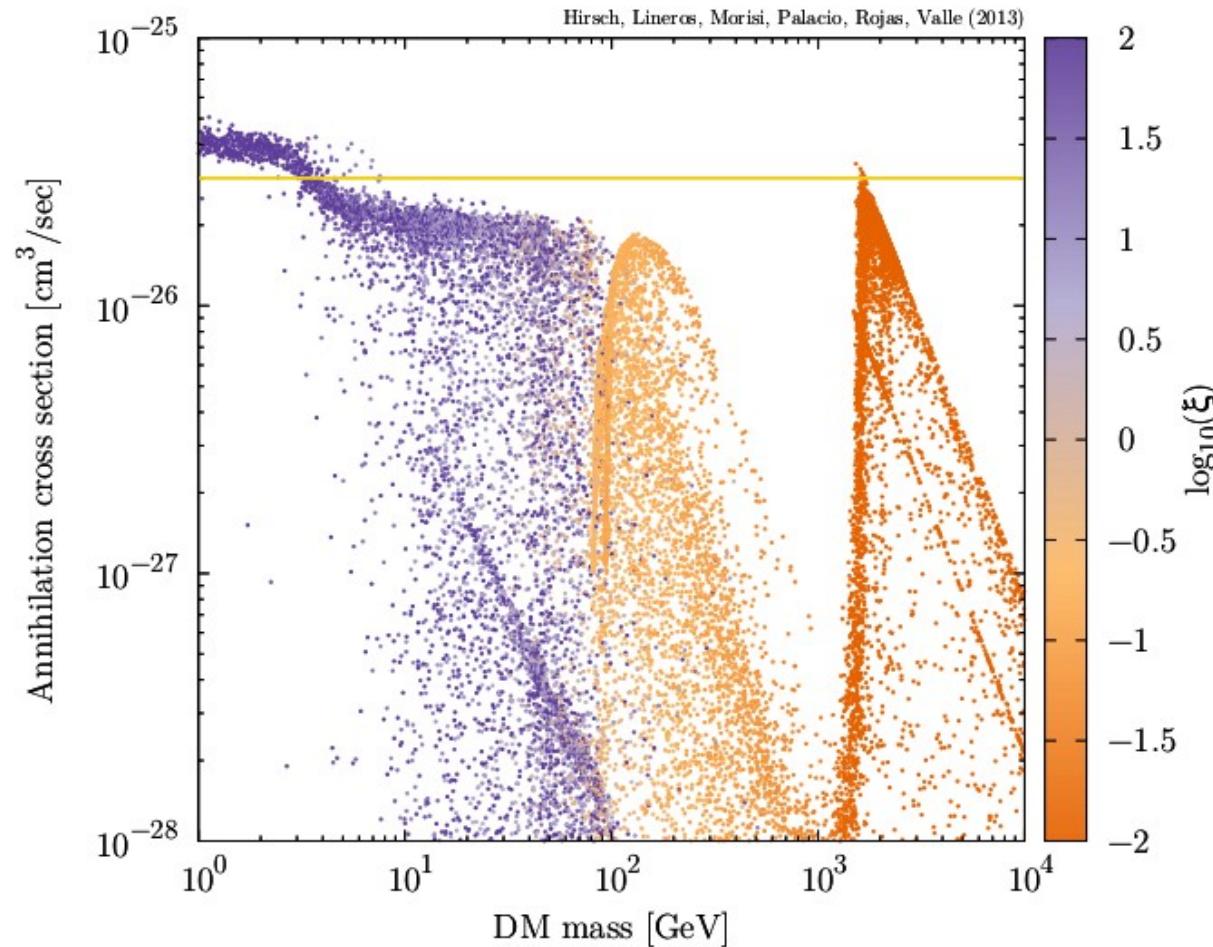
	Standard Model			Fermions		Scalars	
	L	e	ϕ	Σ	N	η	Ω
$SU(2)_L$	2	1	2	3	1	2	3
Y	-1	-2	1	0	0	1	0
Z_2	+	+	+	-	-	-	+

Singlet-Triplet Scotogenic DM

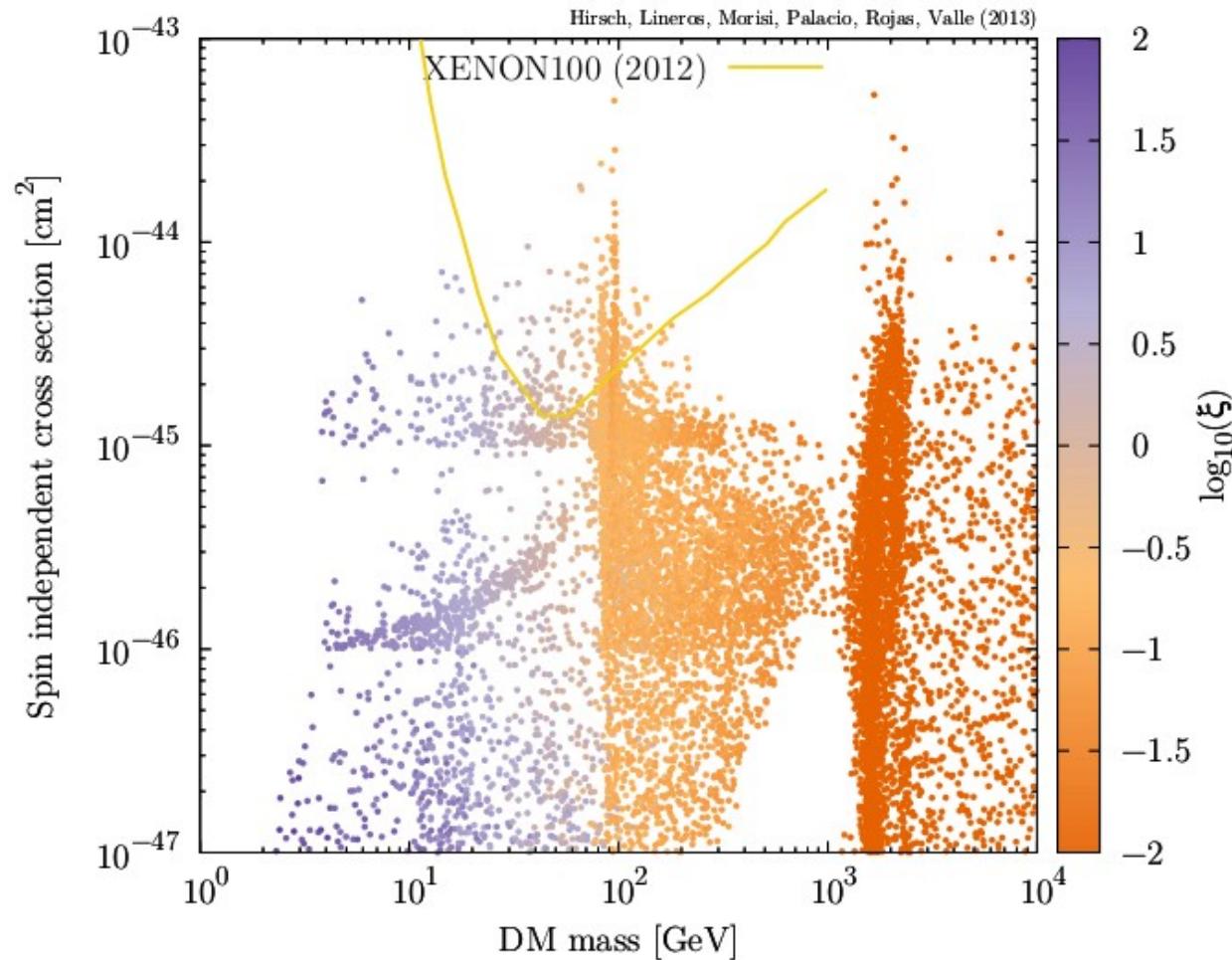
$$\begin{aligned}\mathcal{L} \supset & -Y_{\alpha\beta} \bar{L}_\alpha e_\beta \phi - Y_{\Sigma_\alpha} \bar{L}_\alpha C \Sigma^\dagger \tilde{\eta} - \frac{1}{4} M_\Sigma \text{Tr} [\bar{\Sigma}^c \Sigma] + \\ & -Y_\Omega \text{Tr} [\bar{\Sigma} \Omega] N - Y_{N_\alpha} \bar{L}_\alpha \tilde{\eta} N - \frac{1}{2} M_N \bar{N}^c N + h.c.\end{aligned}$$

$$\begin{aligned}V_{\text{scal}} = & -m_1^2 \phi^\dagger \phi + m_2^2 \eta^\dagger \eta + \frac{\lambda_1}{2} (\phi^\dagger \phi)^2 + \frac{\lambda_2}{2} (\eta^\dagger \eta)^2 + \lambda_3 (\phi^\dagger \phi) (\eta^\dagger \eta) \\ & + \lambda_4 (\phi^\dagger \eta) (\eta^\dagger \phi) + \frac{\lambda_5}{2} (\phi^\dagger \eta)^2 + h.c. - \frac{M_\Omega^2}{4} \text{Tr} (\Omega^\dagger \Omega) + (\mu_1 \phi^\dagger \Omega \phi + h.c.) \\ & + \lambda_1^\Omega \phi^\dagger \phi \text{Tr} (\Omega^\dagger \Omega) + \lambda_2^\Omega (\text{Tr} (\Omega^\dagger \Omega))^2 + \lambda_3^\Omega \text{Tr} ((\Omega^\dagger \Omega)^2) + \lambda_4^\Omega (\phi^\dagger \Omega) (\Omega^\dagger \phi) \\ & + (\mu_2 \eta^\dagger \Omega \eta + h.c.) + \lambda_1^\eta \eta^\dagger \eta \text{Tr} (\Omega^\dagger \Omega) + \lambda_4^\eta (\eta^\dagger \Omega) (\Omega^\dagger \eta),\end{aligned}$$

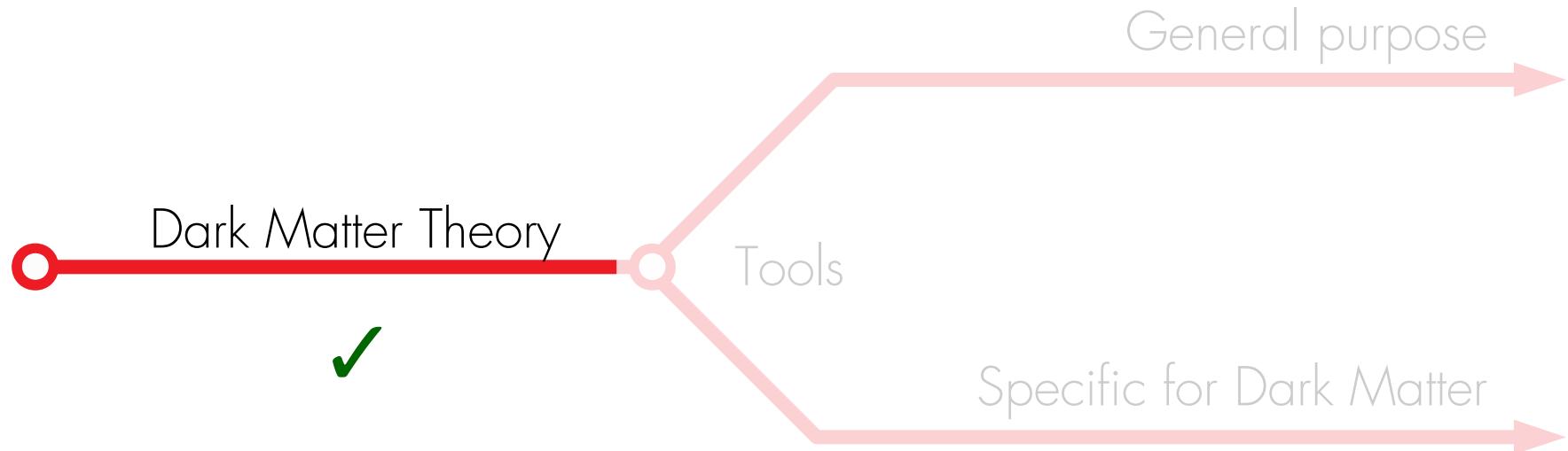
Singlet-Triplet Scotogenic DM

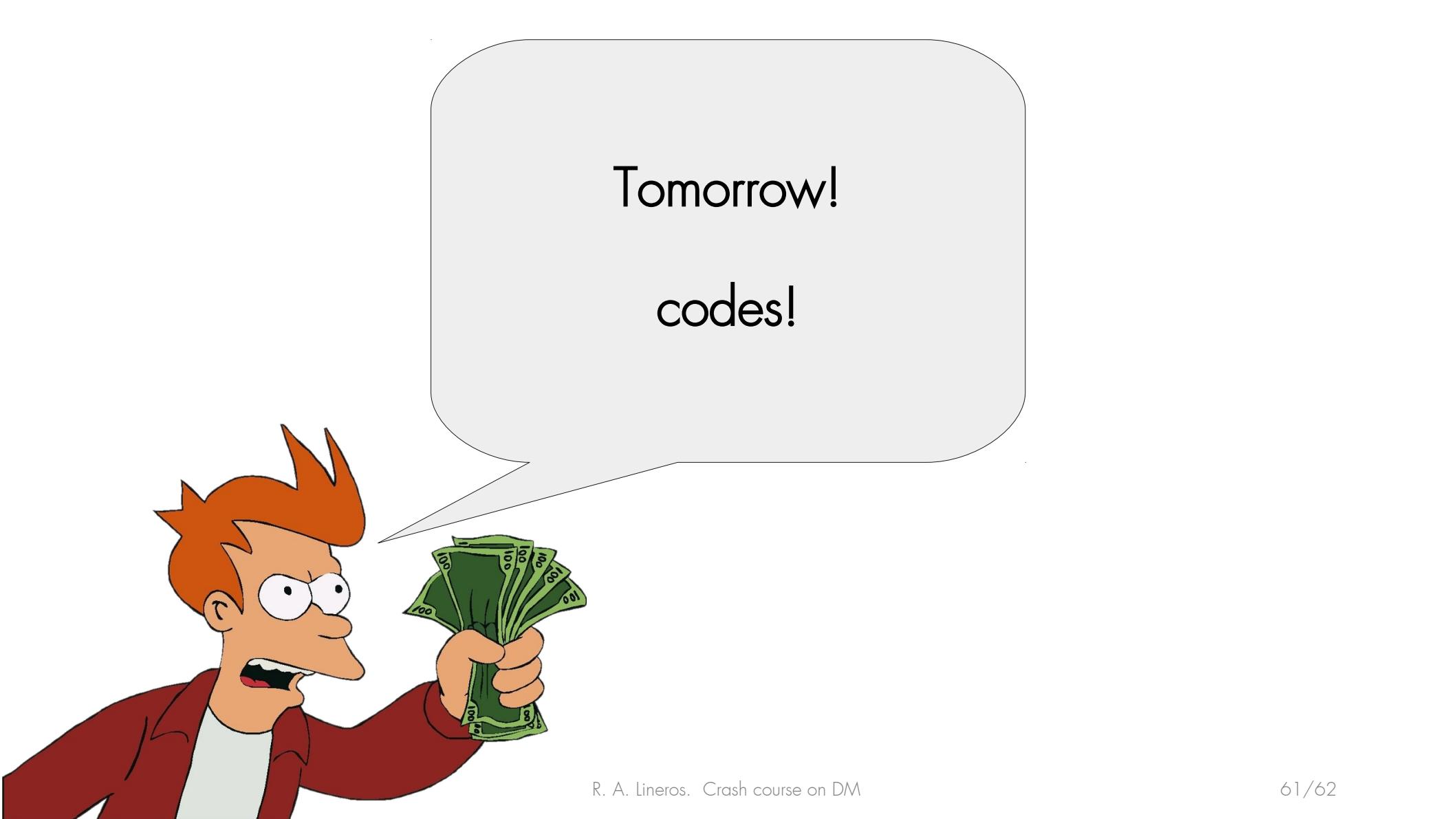


Singlet-Triplet Scotogenic DM



Course plan





Tomorrow!

codes!

Thanks